Advanced Field Theory: Exercise Sheet 6

Due: Thursday, 3.4.2014

Exercise 8 The decay width $\Gamma(\tau^+ \to \overline{\nu}_{\tau} \pi^+)$ can be obtained from the smileptonic pion decay $\pi^+(p) \to \mu^+(k)\nu_{\mu}(q)$ considered in Ex. 7.

a) Start from

$$\sum_{spins} |\mathcal{M}_{\pi^+ \to \mu^+ \nu_{\mu}}|^2 = 8G_F^2 f_{\pi}^2 \left(2(q \cdot p)(k \cdot p) - p^2(q \cdot k) \right) \tag{1}$$

and cross the lepton to the final state and the pion to the initial state via $k \to -k$ and $p \to -p$. Note that there is an additional overall factor of -1, as we have to cross a fermion line. Also, the τ is now in the initial state and hence we should average over the spins. Show that after exploiting the kinematics of the process, we can write

$$\frac{1}{2} \sum_{spins} |\mathcal{M}|^2 = 2G_F^2 f_\pi^2 m_\ell^4 \left(1 - \frac{m_\pi^2}{m_\ell^2} \right).$$
(2)

b) Show that after integration over the phasespace, we obtain

$$\Gamma\left(\tau^{+} \to \overline{\nu}_{\tau} \pi^{+}\right) = \frac{1}{8\pi} G_{F}^{2} f_{\pi}^{2} m_{\ell}^{3} \left(1 - \frac{m_{\pi}^{2}}{m_{\ell}^{2}}\right)^{2}.$$
(3)

Exercise 9 The ρ meson is an isospin triplet of massive spin 1 particles. To compute $\Gamma(\tau^+ \to \overline{\nu}_{\tau} \rho^+)$, we parametrize the matrix element of the vector current between the vacuum and the rho meson as

$$\langle 0|J^{\mu a}|\rho^b_\lambda(p)\rangle = g_\rho \varepsilon^\mu_\lambda e^{-ipx},\tag{4}$$

where g_{ρ} is the rho decay constant, a and b are isospin indices and λ and p are the polarisation index and momentum of the rho meson. The matrix element can be written as

$$\mathcal{M} = 2G_F g_\rho \overline{v}(k) \not \in_{\lambda}(p) P_L v(q).$$
⁽⁵⁾

a) Show that

$$\langle |\mathcal{M}|^2 \rangle = 2G_F^2 g_\rho^2 \mathrm{tr} \left((\not k - m_\ell) \gamma^\mu P_L \not q \gamma^\nu \left(-g_{\mu\nu} + \frac{p_\mu p_\nu}{m_\rho^2} \right) \right). \tag{6}$$

Hint: The polarisation sum for a massive vector particle is

$$\sum_{\lambda} \varepsilon^{\mu}_{\lambda}(p) \varepsilon^{\nu*}_{\lambda}(p) = -g^{\mu\nu} + \frac{p^{\mu}p^{\nu}}{m^2}.$$
(7)

b) Show that only two terms in the trace actually contribute and we can write

$$\langle |\mathcal{M}|^2 \rangle = 2G_F^2 g_\rho^2 \mathrm{tr} \left(-\frac{1}{2} \not{k} \gamma^\mu \not{q} \gamma_\mu + \frac{1}{2m_\rho^2} \not{k} \not{p} \not{q} \not{p} \right).$$
(8)

c) Evaluate the trace to arrive at

$$\langle |\mathcal{M}|^2 \rangle = 2G_F^2 g_\rho^2 \left(2(k \cdot q) + \frac{4(p \cdot q)(k \cdot p)}{m_\rho^2} \right) \tag{9}$$

$$=2G_F^2 g_\rho^2 \frac{m_\ell^4}{m_\rho^2} \left(1 - \frac{m_\rho^2}{m_\ell^2}\right) \left(1 + 2\frac{m_\rho^2}{m_\ell^2}\right),\tag{10}$$

i.e.

$$\Gamma(\tau^+ \to \bar{\nu}_\tau \rho^+) = \frac{1}{8\pi} G_F^2 g_\rho^2 \frac{m_\ell^3}{m_\rho^2} \left(1 - \frac{m_\rho^2}{m_\ell^2} \right)^2 \left(1 + 2\frac{m_\rho^2}{m_\ell^2} \right).$$
(11)