

1) The Lorentz group

The Lorentz group is the set of transformations a four-dimensional space time coordinates (x^μ) , which preserve the scalar product $x \cdot y = x^\mu g_{\mu\nu} y^\nu$

$$x^\mu \rightarrow (x')^\mu = \Lambda^\mu_\nu x^\nu$$
$$y^\mu \rightarrow (y')^\mu = \Lambda^\mu_\nu y^\nu$$

$$g_{\mu\nu} = \text{diag}(1, -1, -1, -1)$$

$$\Lambda^\mu_\lambda \Lambda^\nu_\rho g_{\mu\nu} = g_{\lambda\rho}$$

$$\Rightarrow \boxed{\Lambda^T g \Lambda = g}$$

- * Such transformations form a (non-compact) group
- * The sub-set defined by $\det(\Lambda) = +1, \Lambda^0_0 \geq 1$ is denoted the proper Lorentz group

- * 6 generators: $\begin{cases} J_i \leftrightarrow \text{rotations (3d)} \\ K_i \leftrightarrow \text{boosts} \end{cases}$

$$\Lambda^\mu_\nu = \left[\exp \left(-\frac{i}{2} \omega^{\rho\sigma} M_{\rho\sigma} \right) \right]^\mu_\nu$$

$$M^\mu_\nu = \begin{pmatrix} 0 & -K_1 & -K_2 & -K_3 \\ K_1 & 0 & J_3 & -J_2 \\ K_2 & -J_3 & 0 & J_1 \\ K_3 & J_2 & -J_1 & 0 \end{pmatrix}$$

$$[J_i, J_j] = i \epsilon_{ijk} J_k$$

$$[J_i, K_j] = i \epsilon_{ijk} K_k$$

$$[K_i, K_j] = -i \epsilon_{ijk} J_k$$

- * We can re-arrange the generators into

$$J_i^\pm = \frac{1}{2} (J_i \pm iK_i) \Rightarrow \begin{cases} [J_i^\pm, J_j^\pm] = i \epsilon_{ijk} J_k^\pm \\ [J_i^\pm, J_j^\mp] = \phi \end{cases}$$

⇒ The Lorentz Group is "similar" to a "pair" of $SU(2)$ groups

More precisely $SO(4) \sim SU(2) \times SU(2)$

$SO(1,3) \sim SU(2) \times SU(2)^*$

↑
Lorentz

↑

- ⊙ charge conjugation interchange the two $SU(2)$'s ($(J_i^\pm)^* = -J_i^\mp$)
- ⊙ the representations are not real (as in $so(4)$)

This fact allows us to establish a further group equivalence, namely

$SO(1,3) \sim SL(2, \mathbb{C})$

↑
Group of 2×2 matrices with unit determinant

Actually there is an homomorphism between $SO(1,3)$ & $SL(2, \mathbb{C})$

⇒ for any $A, B \in SL(2, \mathbb{C})$, for any representation of the Lorentz group, there exists $\Lambda(A)$:

$\Lambda(A) \Lambda(B) = \Lambda(A \cdot B)$

Proof: let's consider the matrices

$\sigma^m = (\sigma_0, -\sigma_i)$

$\sigma_0 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$

$\sigma_i =$ Pauli matrices

A) Let's consider the following mapping

$$g: X^\mu \rightarrow X = x_\mu \cdot \sigma^\mu \quad (g: M_4 \rightarrow H)$$

↑
general complex 2x2 matrix that
span the full set of 2x2 Hermitian
complex matrix.
($\sigma^\mu =$ complete set)

$$\det X = \det (x_\mu \sigma^\mu x_\nu \sigma^\nu) = x_\mu x_\nu \det (\sigma^\mu \sigma^\nu) = x_\mu x^\mu$$

\parallel
 $g^{\mu\nu}$

B) Let's consider the following additional mapping

$$a: X \rightarrow X' = A X A^\dagger$$

↑
 $SL(2, \mathbb{C})$

$$\det(X) = \det(X') \quad \& \quad (X')^\dagger = X' \quad \Rightarrow \quad H \xrightarrow{A} H$$

c) Let's finally consider

$$g^{-1}: Y \rightarrow y^\mu = \frac{1}{2} \text{Tr}(Y \cdot \bar{\sigma}^\mu) \quad (g^{-1}: H \rightarrow M_4)$$

↑ same as σ^μ

$$g^{-1}: X = x_\mu \sigma^\mu \rightarrow \frac{1}{2} \text{Tr}(x_\nu \sigma^\nu \bar{\sigma}^\mu) = \frac{1}{2} x_\nu \text{Tr}(\sigma^\nu \bar{\sigma}^\mu) = x^\mu$$

\parallel
 $2 g^{\nu\mu}$



$$M_4 \rightarrow H \rightarrow H \rightarrow M_4$$

$$x_\nu \rightarrow x_\nu \sigma^\nu \rightarrow A x_\nu \sigma^\nu A^\dagger \rightarrow \frac{1}{2} \text{tr}(A x_\nu \sigma^\nu A^\dagger \bar{\sigma}^\mu) = x'^\mu$$



$$\Lambda^\mu_\nu(A) = \frac{1}{2} \text{Tr}(\bar{\sigma}^\mu A \sigma_\nu A^\dagger)$$

The relation $\Lambda^\mu{}_\nu(M)$ $M \in SL(2, \mathbb{C})$
can be inverted up to a sign ambiguity

$$M(\Lambda) = \pm \frac{1}{[\det(\Lambda^\mu{}_\nu \sigma_\mu \bar{\sigma}^\nu)]^{1/2}} \Lambda^\mu{}_\nu \sigma_\mu \bar{\sigma}^\nu$$

$$\Lambda \rightarrow \pm M$$

This correspondence defines a 2-value representation of the Lorentz group that is nothing but the spinor representation

Technically

$$SO(1,3) \Big|_{\substack{\text{proper} \\ \text{transf.}}}$$

$$\leftrightarrow SL(2, \mathbb{C}) / \mathbb{Z}_2$$