

Spontaneous SUSY breaking

If SUSY is spontaneously broken, it means that the ground state is not invariant under SUSY transformations

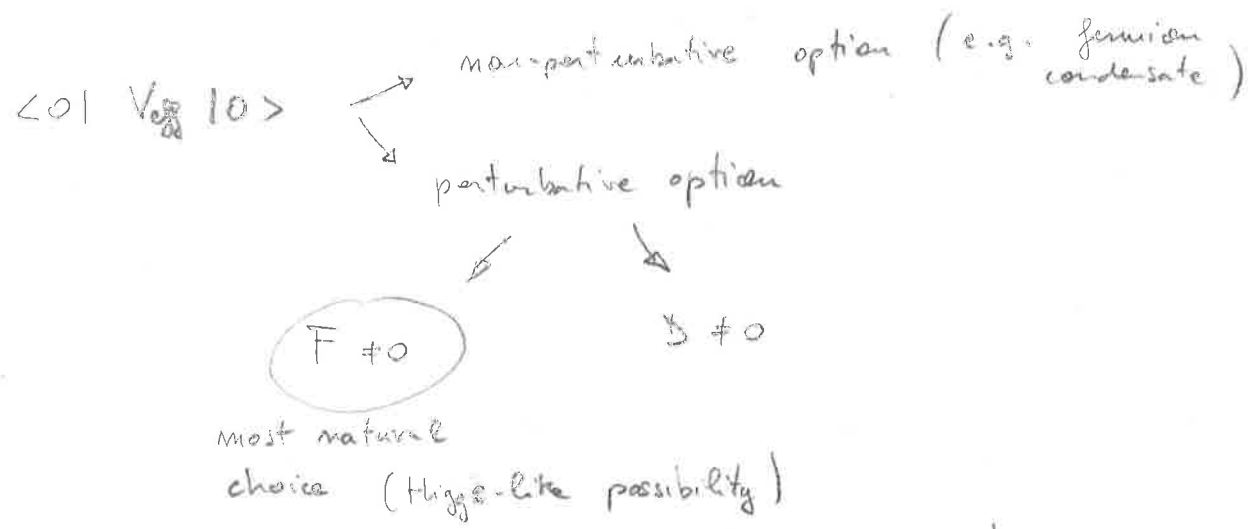
$$\Rightarrow Q_i |0\rangle \neq 0 \quad \& \quad \bar{Q}_i |0\rangle \neq 0$$

↑
ground state

Since $H = P^0 = \frac{1}{4} (Q_1 \bar{Q}_1 + Q_2 \bar{Q}_2 + \bar{Q}_1 Q_1 + \bar{Q}_2 Q_2)$

this means that $\langle 0 | H | 0 \rangle = \frac{1}{4} \sum_i (|Q_i |0\rangle|^2 + |\bar{Q}_i |0\rangle|^2) > 0$

⇒ we must have an effective potential such that the ground state has positive energy



(Recall $V(\psi_i, \bar{\psi}_i) = \sum_i \left| \frac{\partial W}{\partial \psi_i} \right|^2 = \sum_i F_i \bar{F}_i$)

The O'Raifeartaigh model

$$W = -\Lambda^2 \Phi_1 + m \Phi_1 \Phi_3 + \frac{g}{2} \Phi_1 \Phi_3^2$$

$\Phi_i =$ chiral superfields (not charged under gauge symm.)

N.B.: The linear term in Φ_1 is possible

ONLY if Φ_1 has no charges \Rightarrow not possible

in the MSSM \Rightarrow need for extra fields

(general conclusion for susy breaking)

$$V(\phi_i) = \sum_{i=1}^3 |\mathcal{F}_i|^2$$

$$\mathcal{F}_i = \left. \frac{\partial V}{\partial \Phi_i} \right|_{\phi_i = \phi_i}$$

$$\mathcal{F}_1 = -\Lambda^2 + \frac{y}{2} \phi_3^2$$

$$\mathcal{F}_2 = m \phi_3$$

$$\mathcal{F}_3 = m \phi_2 - y \phi_1 \phi_3$$

It is clear that we cannot simultaneously have $\mathcal{F}_1 = \mathcal{F}_2 = \mathcal{F}_3 = 0$

\Rightarrow good prototype model for susy breaking.

$$V \Big|_{\substack{\phi_2=0 \\ \phi_1=0}} = \left| \Lambda^2 - \frac{y}{2} \phi_3 \right|^2 + m^2 |\phi_3|^2$$

$$\longrightarrow V \geq \Lambda^2$$

assume $m^2 > y \Lambda^2$

(min. configuration obtained for $\phi_2 = \phi_3 = 0$ any ϕ_1)

$$\boxed{\sqrt{V} \Big|_{\text{min}} = \Lambda}$$

Expanding around $\phi_1 = \phi_2 = \phi_3 = 0$ one gets the following

(tree-level) spectrum:

scalar masses

- 0 } masses associate to ϕ_1
- 0 } (flat direction)
- m^2
- m^2
- $m^2 - y \Lambda^2$
- $m^2 + y \Lambda^2$

fermion masses

- 0 \leftarrow GOLDSTINO
- m
- m

Some comments:

- (A) The different mass spectrum between fermions & bosons is a clear signal that SUSY has been broken
- (B) The "flat direction" in ϕ_1 is an accident of the tree-level potential \rightarrow removed by quantum corrections \rightarrow non-vanishing masses for all scalars beyond the tree-level
- (C) The massless fermion, on the other hand, is a general consequence of the spontaneous symmetry breaking of SUSY (analog of Goldstone boson)
- \hookrightarrow it can be removed only gauging SUSY (supergravity) \Rightarrow mass to the GRAVITINO
- (D) Note that $\text{Tr}[M^2_{\text{scalars}}] = 2 \text{Tr}[M^2_{\text{fermions}}]$
- This property also remains valid in general



From (D) + The fact that ϕ_1 cannot be changed we deduce that SUSY breaking must occur in a different "sector" of the theory (new fields) and is "communicated" in some way to the "SM sector"

"Hidden sector"
 (non-SM chiral fields)
 ↓
 SUSY breaking

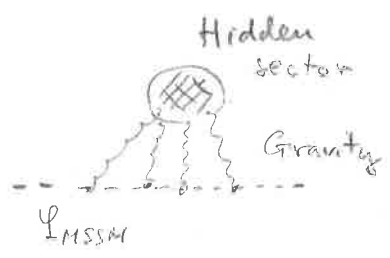
Flavor-blind interactions

MSSM
 "Visible sector"

Two classical way to "mediate" the breaking of SUSY from the hidden sector to the visible sector is via

- ⊙ gravitational interactions (gravity-mediated MSSM)
- ⊙ gauge interactions (gauge-mediated MSSM)

Gravity mediation is always present, by naive dim. analysis we should expect

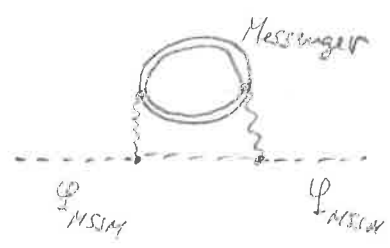


$$M_{soft}^2 \Big|_{\text{Grav.}} \sim \frac{(\sqrt{F})^4}{M_{\text{Planck}}^2}$$

$$\rightarrow (\sqrt{F}) \gtrsim 10^{10} \text{ GeV} \Rightarrow M_{soft} \gtrsim \text{few } \times (100 \text{ GeV})$$

Gauge mediation is possible only if some of the particles in the hidden sector (called "messengers") have non-zero gauge interactions:

$$M_{soft}^2 \sim \left(\frac{g_{SM}}{16\pi^2} \right)^2 \frac{(\sqrt{F})^4}{M_{\text{mess}}^2}$$



mass of the messenger

$$M_{soft} \sim \frac{\alpha_{SM}}{4\pi} \frac{(\sqrt{F})^2}{M_{\text{mess}}} \rightarrow \frac{\alpha_{SM}}{4\pi} \sqrt{F} \Rightarrow \sqrt{F} \gtrsim 10^6 \text{ GeV}$$

(if $M_{\text{mess}} \sim \sqrt{F}$) $\rightarrow M_{soft} \gtrsim 1 \text{ TeV}$

The key point is that in both cases we expect flavor-degenerate soft-breaking terms (up to corrections induced by the Yukawa couplings) \rightarrow no problem with flavor.

The simplest scenario proposed in the literature is the so-called minimal/constrained gravity mediation framework, where one assumes (in an EFT-type inspired approach):

$$\mathcal{L}_{\text{soft}}^{\text{mSUGRA}} = -\frac{1}{M_{\text{Pl}}} F_S \sum_a \frac{1}{2} g_a \lambda^a \lambda^a + \text{h.c.}$$

$$- \frac{1}{M_{\text{Pl}}^2} |F_S|^2 \sum_i g_{\phi} (\phi_i \phi_i^*)$$



$$M_3 = M_2 = M_1 = g \frac{F_S}{M_{\text{Pl}}} \stackrel{\text{def}}{=} m_{1/2}$$

$$M_Q^2 = m_{u_c}^2 = m_{d_c}^2 = m_L^2 = m_{e_c}^2 = m_0^2 \cdot 1/3 \quad m_0^2 \stackrel{\text{def}}{=} g_{\phi} \frac{F_S^2}{M_{\text{Pl}}^2}$$

and similarly for the A terms and the soft Higgs mass term (\rightarrow only 4 parameters!)

But many other options are possible \Rightarrow data will tell (if SUSY will be discovered...)