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Issued: 23.05.2014
Due: 30.05.2014

Exercise 1 [Post-Newtonian Lagrangian for a binary system] (10 points)

The gravitational action for a particle a propagating in the metric $g_{\mu\nu}$ is given by

$$S_a = -m_a c \int dt \left(-g_{\mu\nu} \frac{dx_a^\mu}{dt} \frac{dx_a^\nu}{dt} \right)^{1/2}, \quad (1)$$

where $t = x^0/c$ (not the proper time of the particle) and $dx^\mu/dt = (-c, \mathbf{v})$. In this problem we are considering two masses m_1 and m_2 with separation r . Note that the total action is given by $S = S_1 + S_2$ (paying attention to not double count the terms of the potential energy) and that particle 1 propagates in the metric generated by particle 2 and vice versa.

- (i) Expand the action to order 1PN. Using the definitions:

$${}^{(2)}g_{00} = -2\phi, \quad {}^{(2)}g_{ij} = -2\phi\delta_{ij}, \quad (2)$$

$${}^{(3)}g_{0i} = \zeta_i + \partial_i\partial_0\chi, \quad {}^{(4)}g_{00} = -2(\phi^2 + \psi), \quad (3)$$

you should finally get, for e.g. particle m_1 :

$$\mathcal{L} = m_1 c^2 \left\{ \frac{1}{2} \left(\frac{v_1}{c} \right)^2 - \phi + \frac{1}{8} \left(\frac{v_1}{c} \right)^4 - \frac{\phi^2}{2} - \psi - \frac{3\phi}{2} \left(\frac{v_1}{c} \right)^2 + \zeta_i \frac{v_1^i}{c} + \frac{v_1^i}{c} \partial_i \partial_0 \chi \right\}. \quad (4)$$

(Hint: remember that 1PN means to expand up to ϵ^4 , where $\epsilon = v/c$).

- (ii) Expand the energy momentum tensor:

$$T^{\mu\nu} = \frac{1}{\sqrt{-g}} \sum_a \frac{dt}{d\tau_a} m_a \frac{dx_a^\mu}{dt} \frac{dx_a^\nu}{dt} \delta^{(3)}(\mathbf{x} - \mathbf{x}_a(t)) \quad (5)$$

to order 1PN. Here τ_a is the proper time of particle a .

(Hint: remember that the metric in which m_1 propagates is given by m_2 only, and so $T^{\mu\nu}$. Moreover, only terms up to ϵ^2 are necessary in the expansion of $T^{\mu\nu}$. Note also that $d\tau/dt = \mathcal{L}$.)

- (iii) From the results at point (ii), compute the actual values of the terms in \mathcal{L} , using the definitions:

$$\phi(t, \mathbf{x}) = -\frac{G}{c^4} \int \frac{d^3x'}{|\mathbf{x} - \mathbf{x}'|} {}^{(0)}T^{00}(t, \mathbf{x}'), \quad (6)$$

$$\zeta_i(t, \mathbf{x}) = -\frac{4G}{c^4} \int \frac{d^3x'}{|\mathbf{x} - \mathbf{x}'|} {}^{(1)}T^{0i}(t, \mathbf{x}'), \quad (7)$$

$$\psi(t, \mathbf{x}) = -\frac{G}{c^4} \int \frac{d^3x'}{|\mathbf{x} - \mathbf{x}'|} \left[{}^{(2)}T^{00}(t, \mathbf{x}') + {}^{(2)}T^{ii}(t, \mathbf{x}') \right], \quad (8)$$

$$\chi(t, \mathbf{x}) = -\frac{G}{2c^4} \int d^3x' |\mathbf{x} - \mathbf{x}'| {}^{(0)}T^{00}(t, \mathbf{x}'). \quad (9)$$

Hint: while computing ψ , there is a term in ${}^{(2)}T^{00}$ that depends on ϕ . For a system of 2 particles only, it contributes as 0 in computing ψ and you can neglect it.)

(iv) Plug the metric perturbations into the Lagrangian. For e.g. particle m_1 , you will get:

$$\mathcal{L} = \frac{1}{2}m_1v_1^2 + \frac{Gm_1m_2}{r_{12}} + \frac{1}{8}m_1\frac{v_1^4}{c^2} - \frac{G^2m_2^2m_1}{r_{12}^2c^2} \quad (10)$$

$$+ \frac{Gm_1m_2}{r_{12}} \left\{ 3 \left[\left(\frac{v_1}{c} \right)^2 + \left(\frac{v_2}{c} \right)^2 \right] - 7 \frac{\mathbf{v}_1 \cdot \mathbf{v}_2}{c^2} - \frac{(\mathbf{v}_1 \cdot \hat{\mathbf{r}}_{12})(\mathbf{v}_2 \cdot \hat{\mathbf{r}}_{12})}{c^2} \right\}. \quad (11)$$

Sum it up with the equivalent lagrangian for particle m_2 (do not forget to symmetrize all the terms in which both index 1 and 2 appear symmetrically, which means counting them only once!). You will get finally:

$$\mathcal{L} = \frac{1}{2}m_1v_1^2 + \frac{1}{2}m_2v_2^2 + \frac{Gm_1m_2}{r_{12}} + \frac{1}{8}m_1\frac{v_1^4}{c^2} + \frac{1}{8}m_2\frac{v_2^4}{c^2} \quad (12)$$

$$- \frac{G^2m_2m_1(m_1 + m_2)}{2r_{12}^2c^2} \quad (13)$$

$$+ \frac{Gm_1m_2}{r_{12}} \left\{ 3 \left[\left(\frac{v_1}{c} \right)^2 + \left(\frac{v_2}{c} \right)^2 \right] - 7 \frac{\mathbf{v}_1 \cdot \mathbf{v}_2}{c^2} - \frac{(\mathbf{v}_1 \cdot \hat{\mathbf{r}}_{12})(\mathbf{v}_2 \cdot \hat{\mathbf{r}}_{12})}{c^2} \right\}. \quad (14)$$