

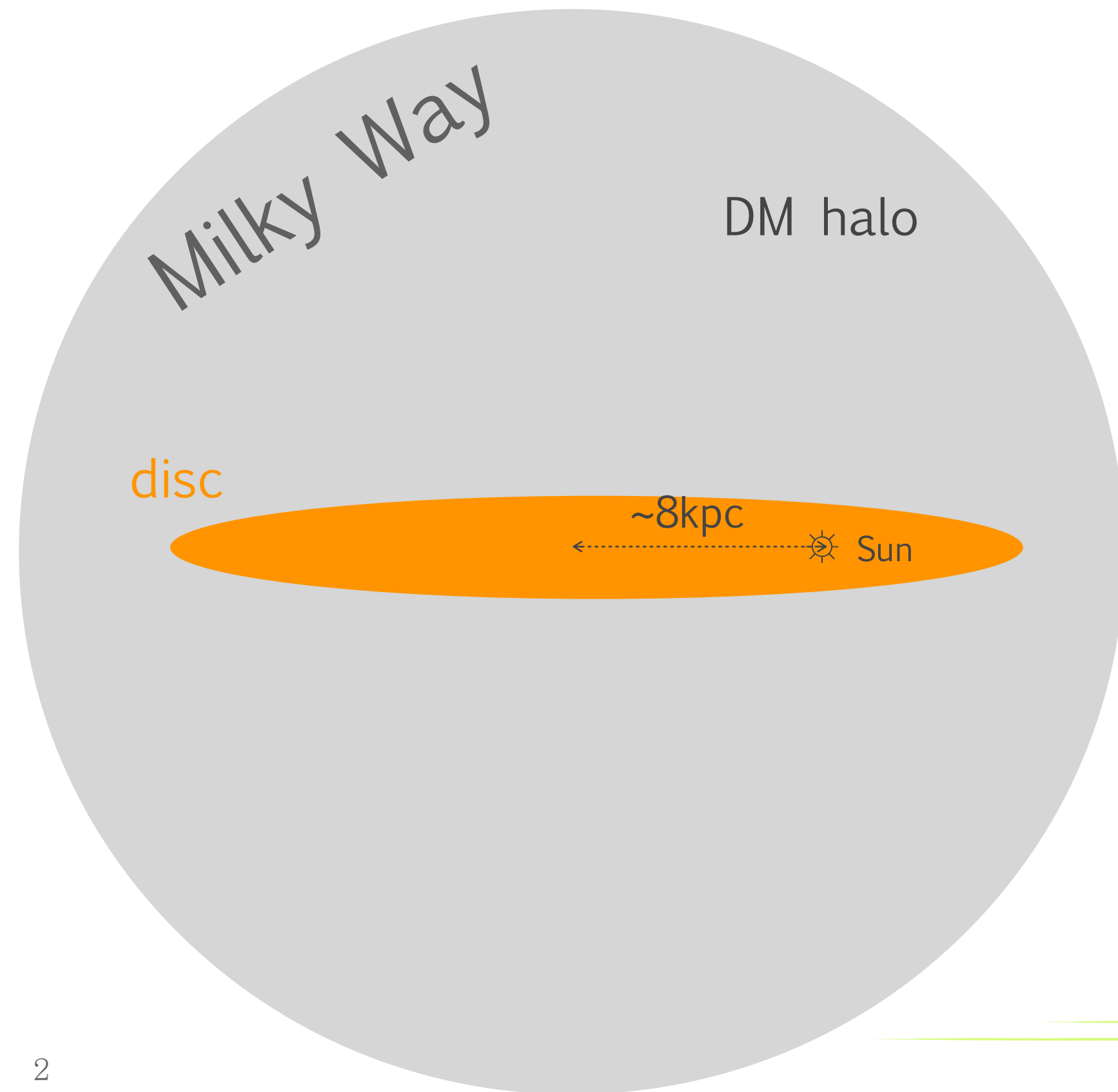


A new measurement of $\rho_{\text{DM}}(R_{\odot})$ from the kinematics of K dwarfs

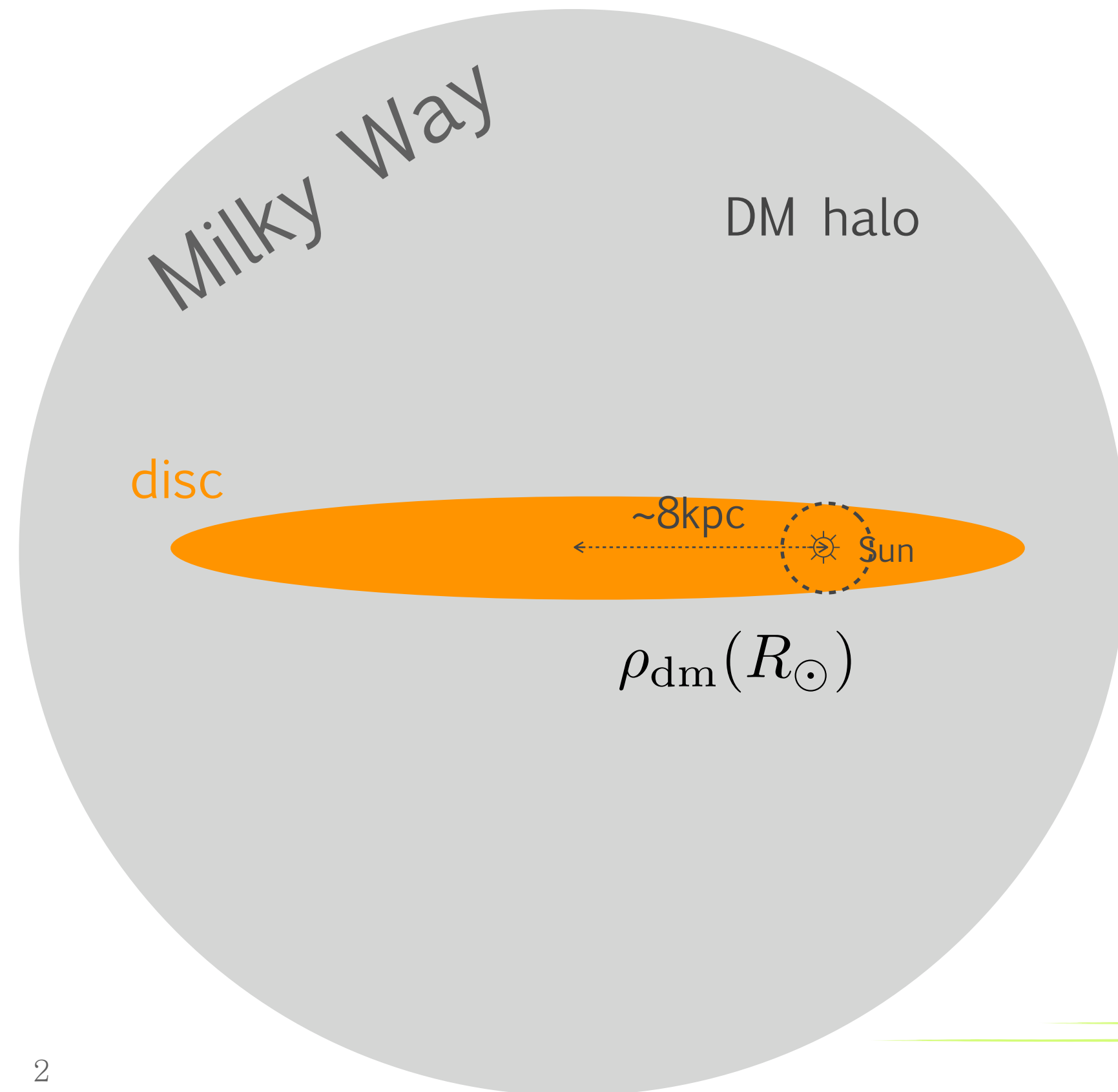
Silvia Garbari

with Justin Read, George Lake and Chao Liu

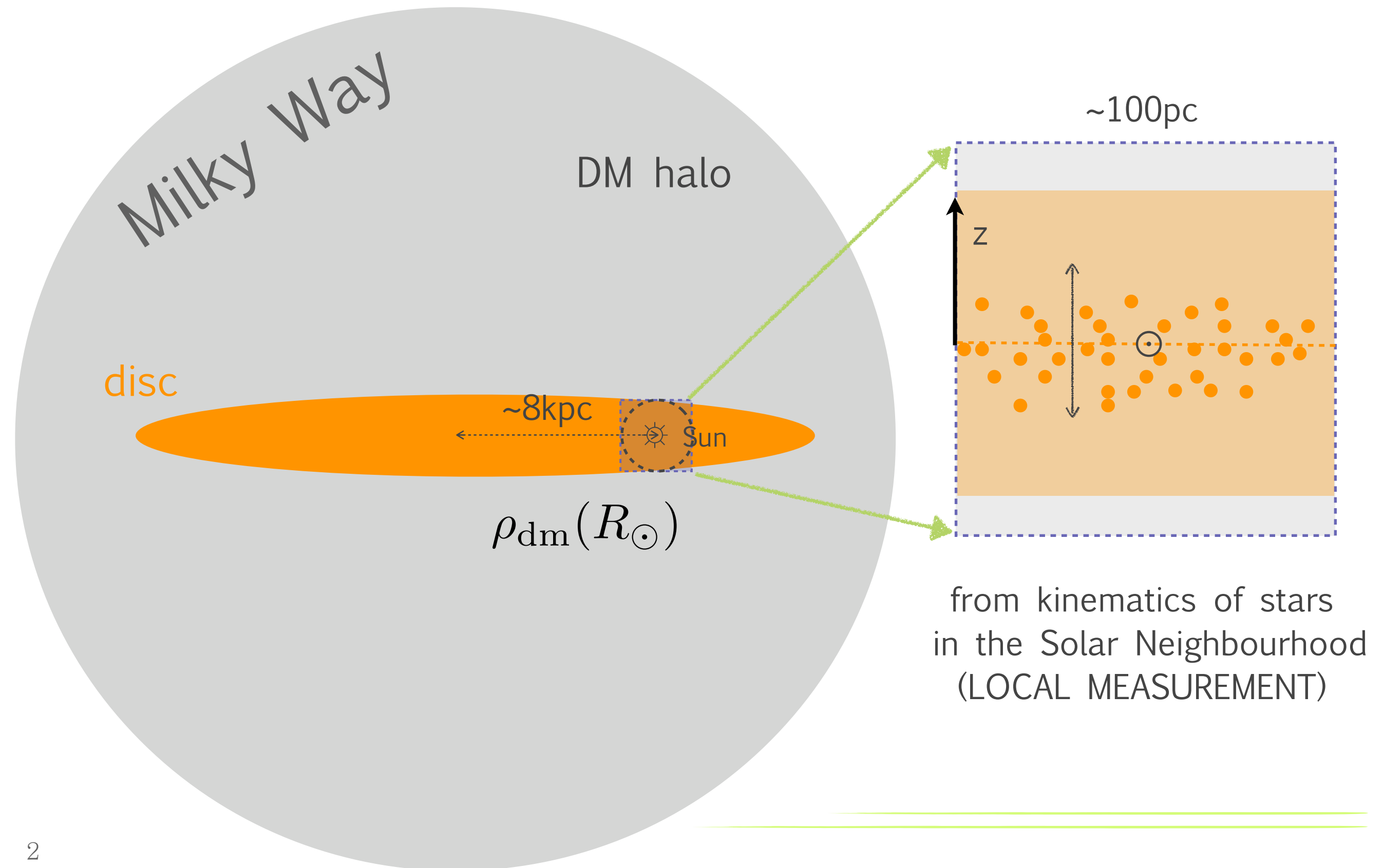
Measuring the local dark matter density $\rho_{\text{DM}}(R_{\odot})$



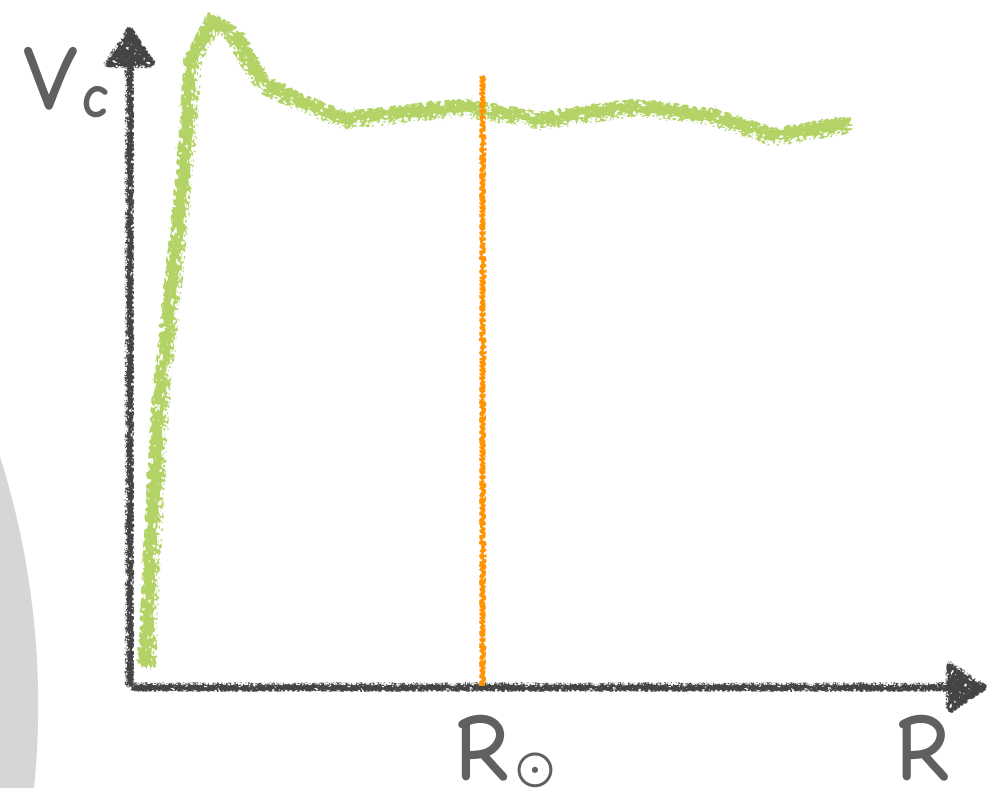
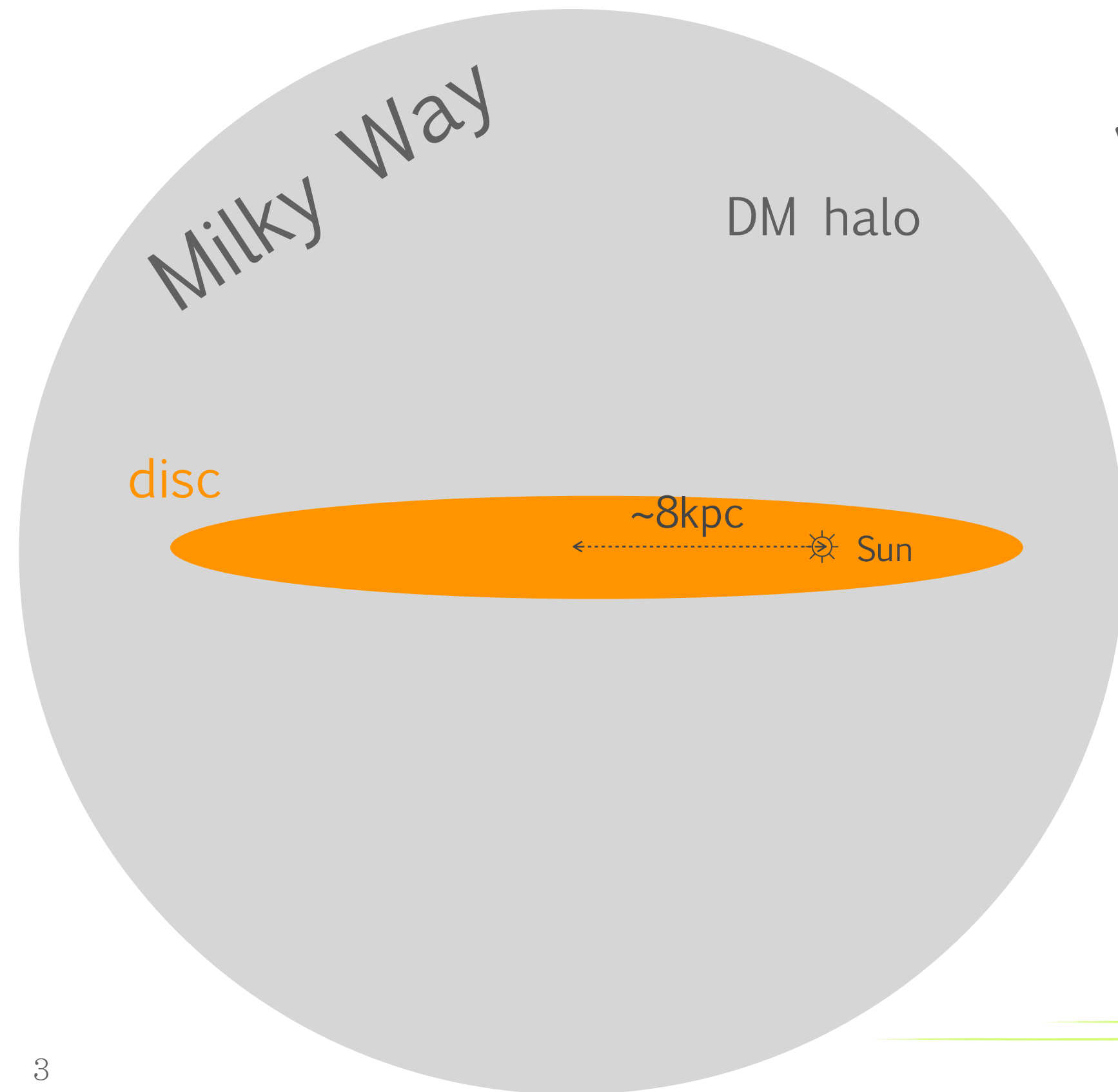
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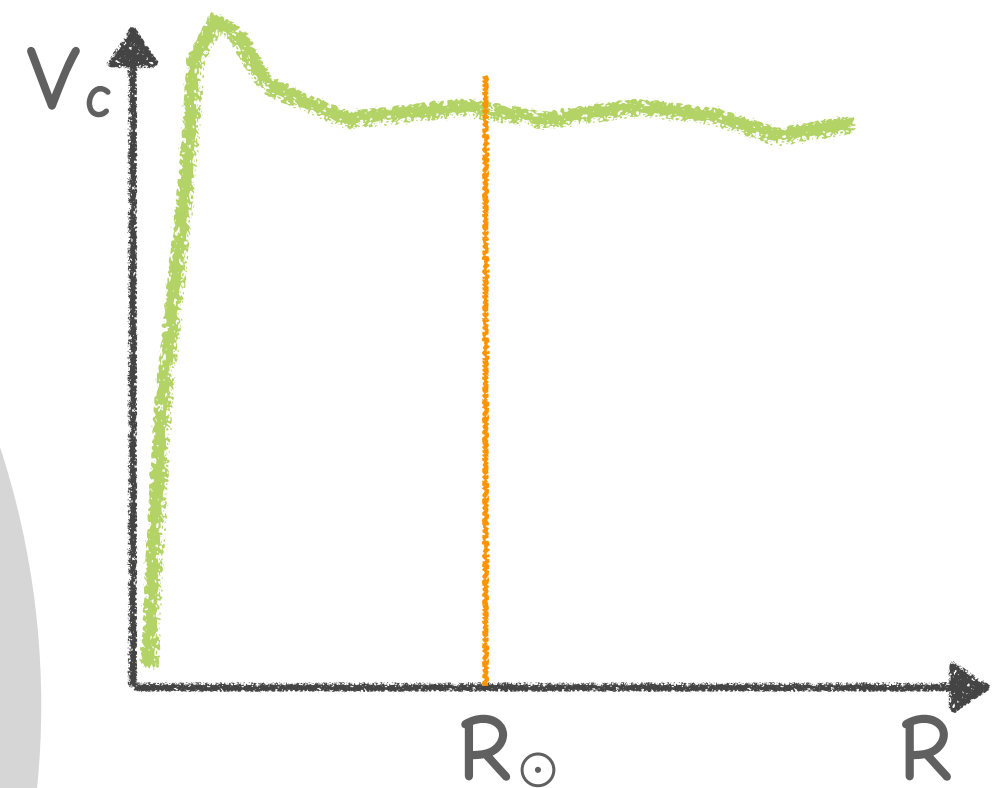
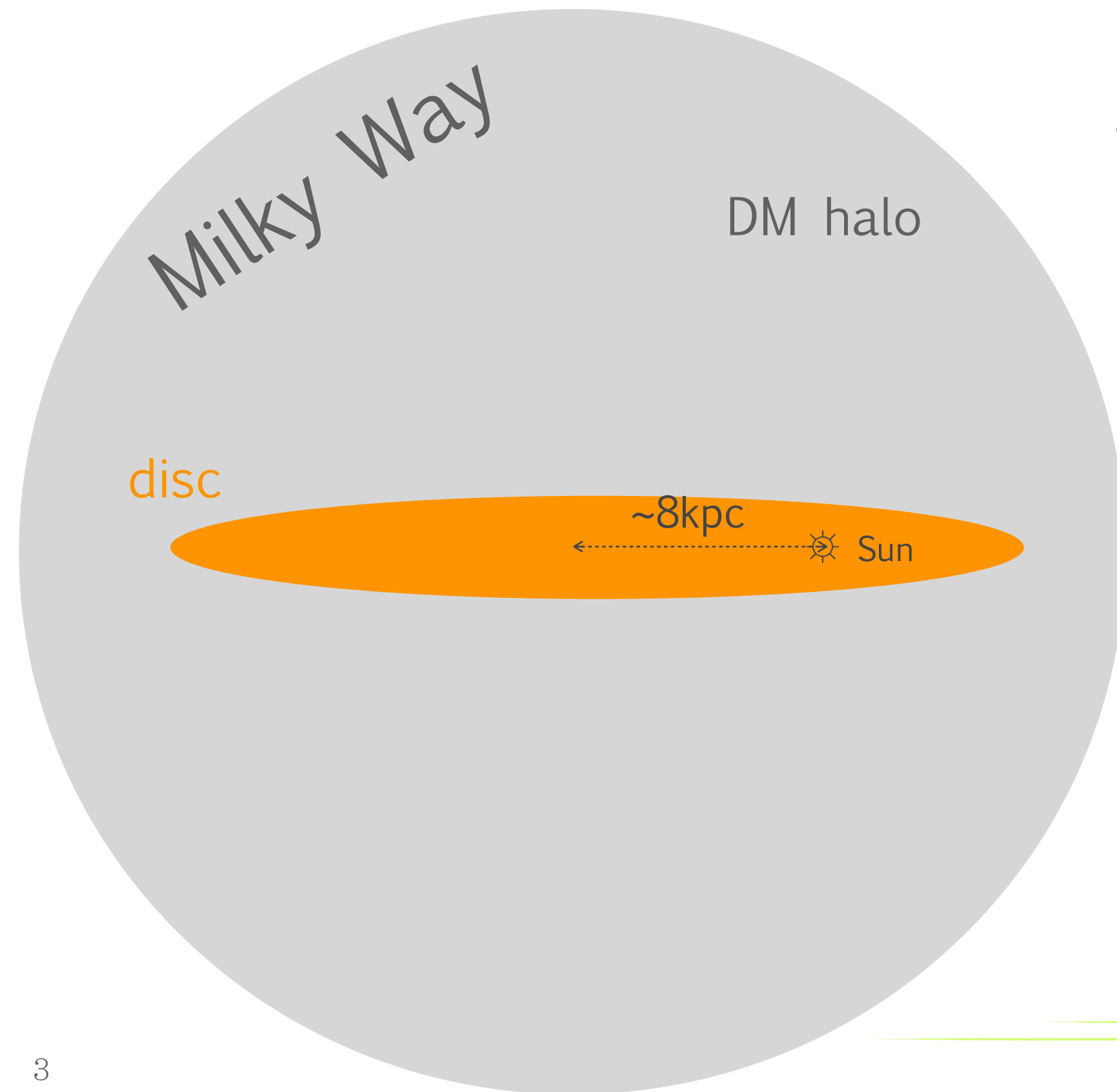
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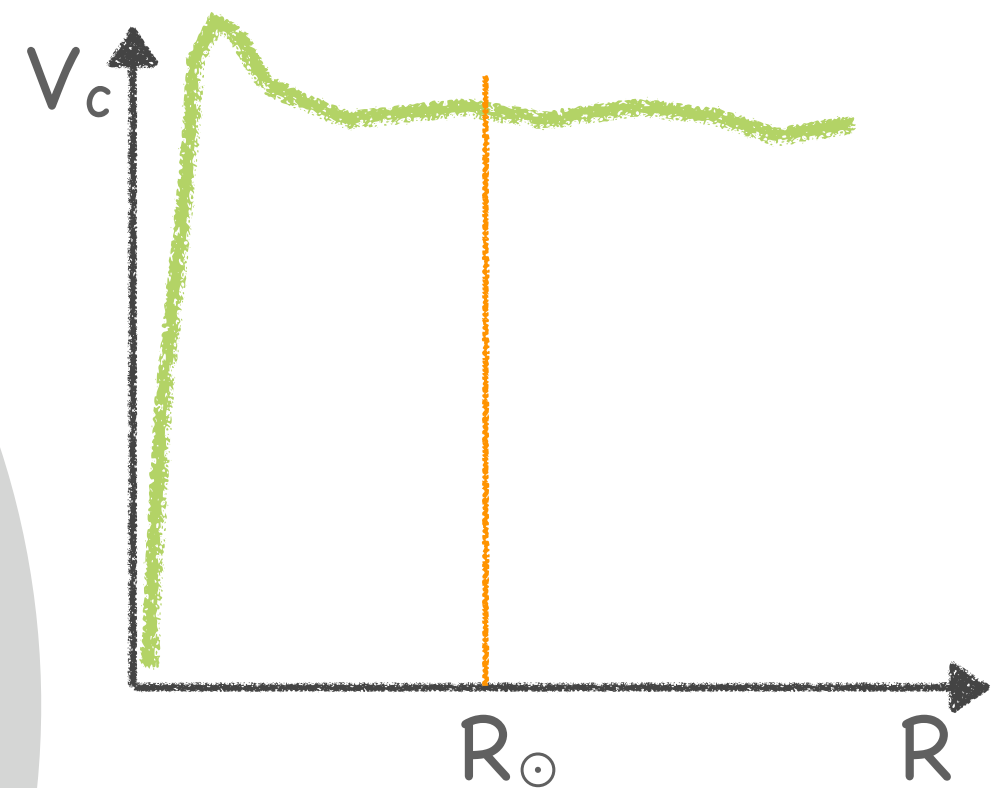
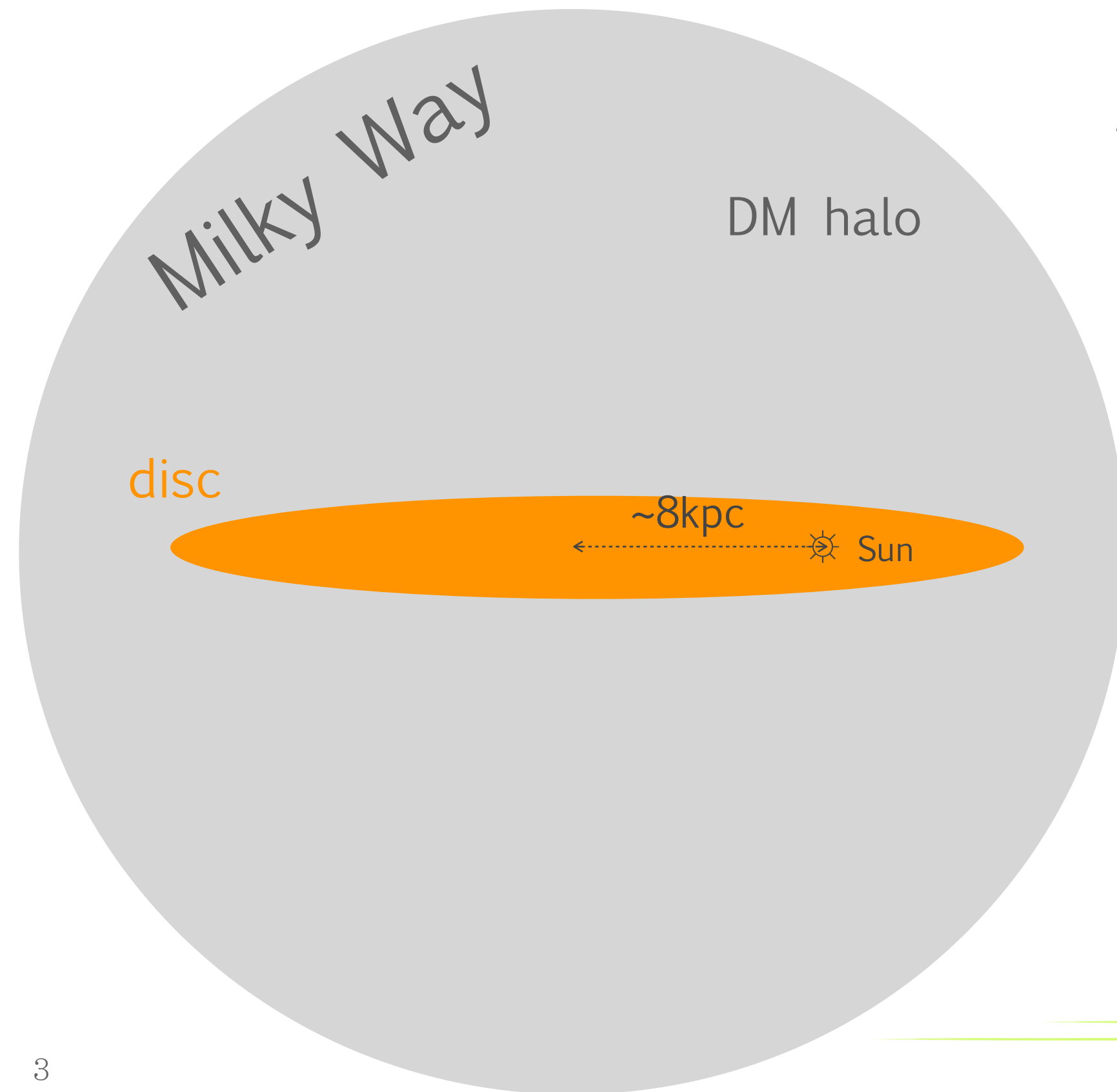


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assumptions
on the halo shape

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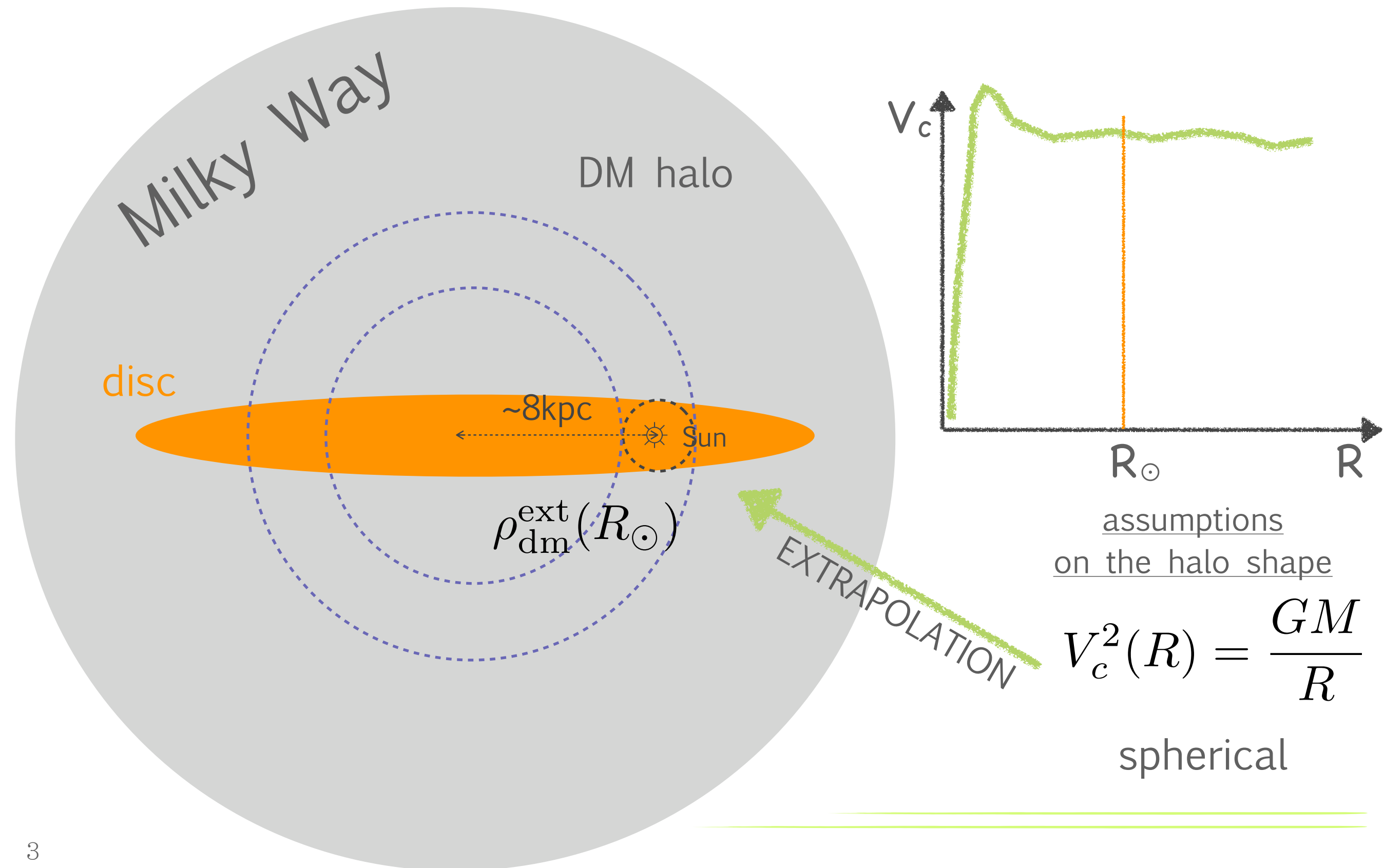


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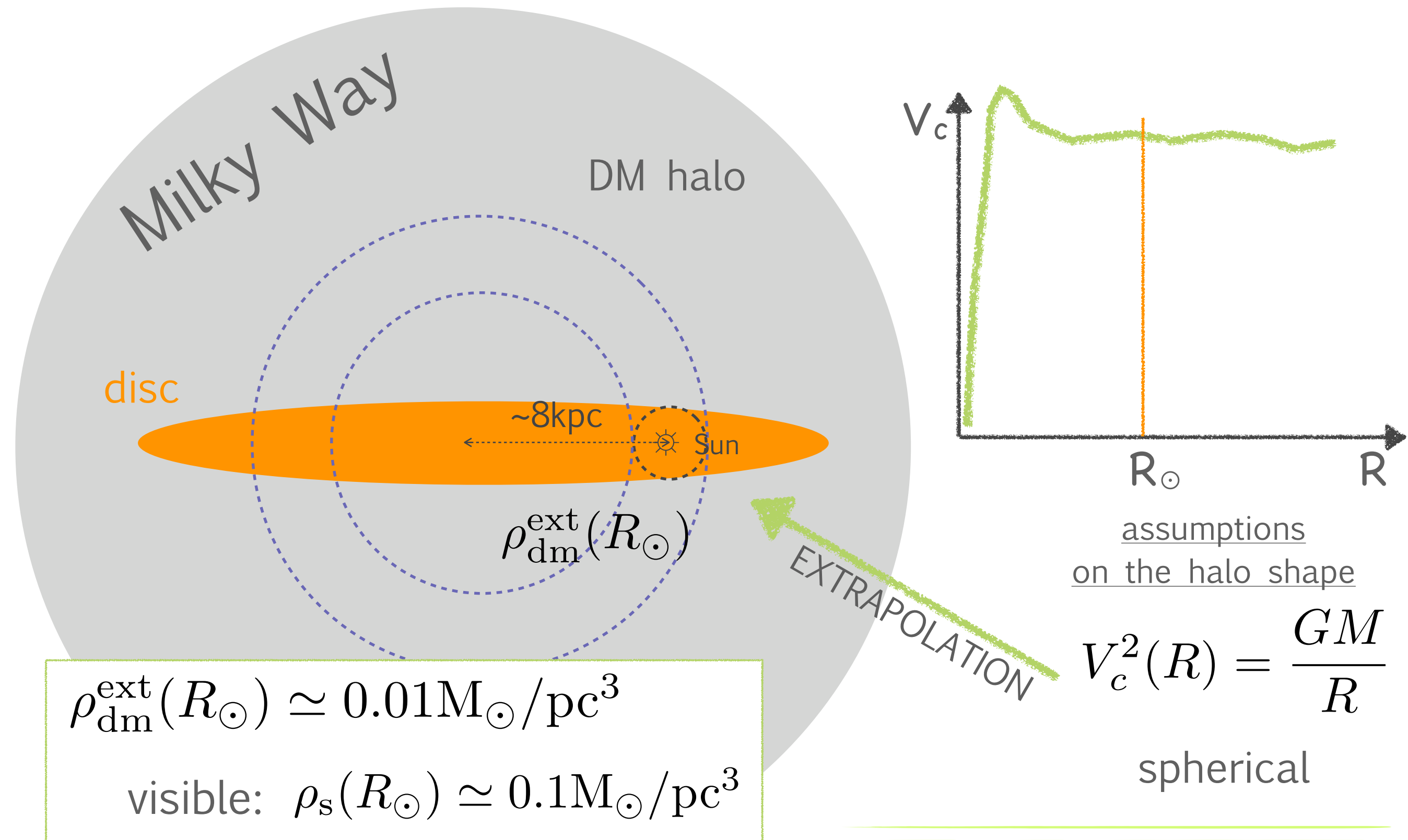
$$V_c^2(R) = \frac{GM}{R}$$

spherical

Measuring the local dark matter density $\rho_{\text{DM}}(R_{\odot})$



Measuring the local dark matter density $\rho_{\text{DM}}(R_{\odot})$



$$\rho_{\text{dm}}^{\text{ext}}(R_{\odot}) \simeq 0.01 M_{\odot} / \text{pc}^3$$

$$\text{visible: } \rho_{\text{s}}(R_{\odot}) \simeq 0.1 M_{\odot} / \text{pc}^3$$

Why is $\rho_{\text{DM}}(R_{\odot})$ interesting?

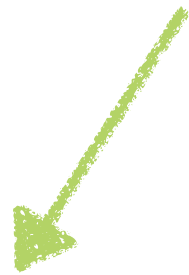
1. Constraints on the shape of the DM Halo of the MW

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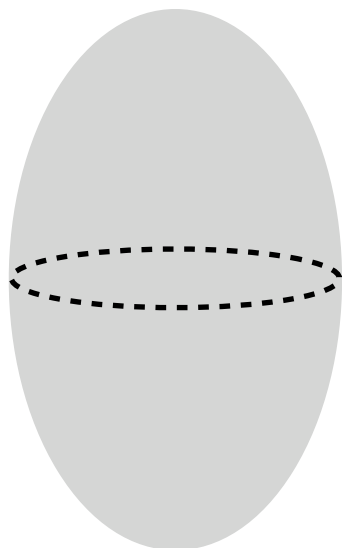
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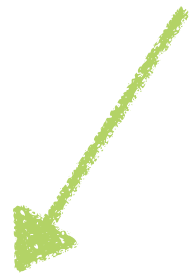
prolate halo



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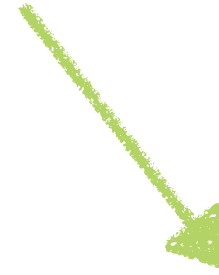
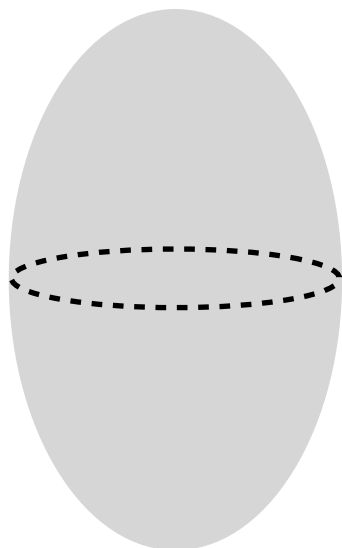
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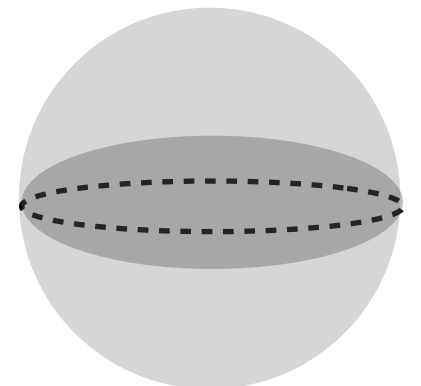
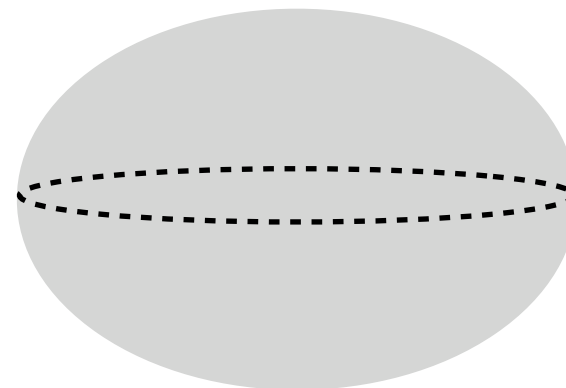
prolate halo



$$\rho_{\text{dm}}(R_{\odot}) > \rho_{\text{dm}}^{\text{ext}}(R_{\odot})$$

oblate halo

(or/and dark disc [Read 2008,2009])



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2. WIMPS detection experiments

if DM=WIMP (Weakly Interacting Massive Particles)

direct detection:

nuclear recoil caused by a WIMP scattering within the detector

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what we want to measure

Minimal Assumption Method

$$\frac{df}{dt} = 0 = \frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{r}} - \nabla \Phi \cdot \frac{\partial f}{\partial \mathbf{v}}$$

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hyp. 1: the system is in equilibrium
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Moments =
Jeans equations

$$\frac{1}{R} \frac{\partial}{\partial R} (R \nu_i \overline{v_R v_z}) + \frac{\partial}{\partial z} \left(\nu_i \overline{v_z^2} \right) + \nu_i \frac{\partial \Phi}{\partial z} = 0$$

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hyp. 2: tilt term is negligible

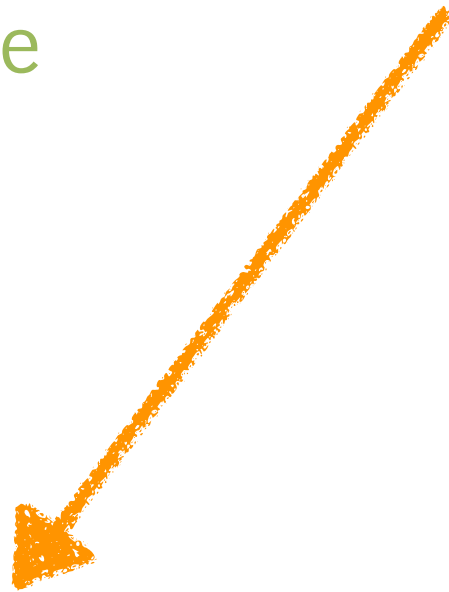
$\Leftrightarrow z_d \ll R_d$ [Bahcall, 1984; Binney&Tremaine 2008]

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Tracer population:

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Poisson: matter density

$$4\pi G(\rho_s + \rho_{\text{dm}}) = \frac{\partial^2 \Phi}{\partial z^2} + \frac{1}{R} \frac{\partial}{\partial R} \left(R \frac{\partial \Phi}{\partial R} \right)$$

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negligible

$$\frac{1}{R} \frac{\partial}{\partial R} \left(R \frac{\partial \Phi}{\partial R} \right) = \frac{1}{R} \frac{\partial V_c^2}{\partial R} = 2(B^2 - A^2)$$

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MCMC

HF method [Holmberg & Flynn, 2000]

hyp. 5: the distribution function is separable

$$f = f_{R,\theta}(v_R, v_\theta, R) \times f_z(v_z, z)$$

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$$\begin{array}{l} f(v_z(0), 0) \\ \Phi(z) \end{array}$$

observed

calculated

Testing the methods with simulations

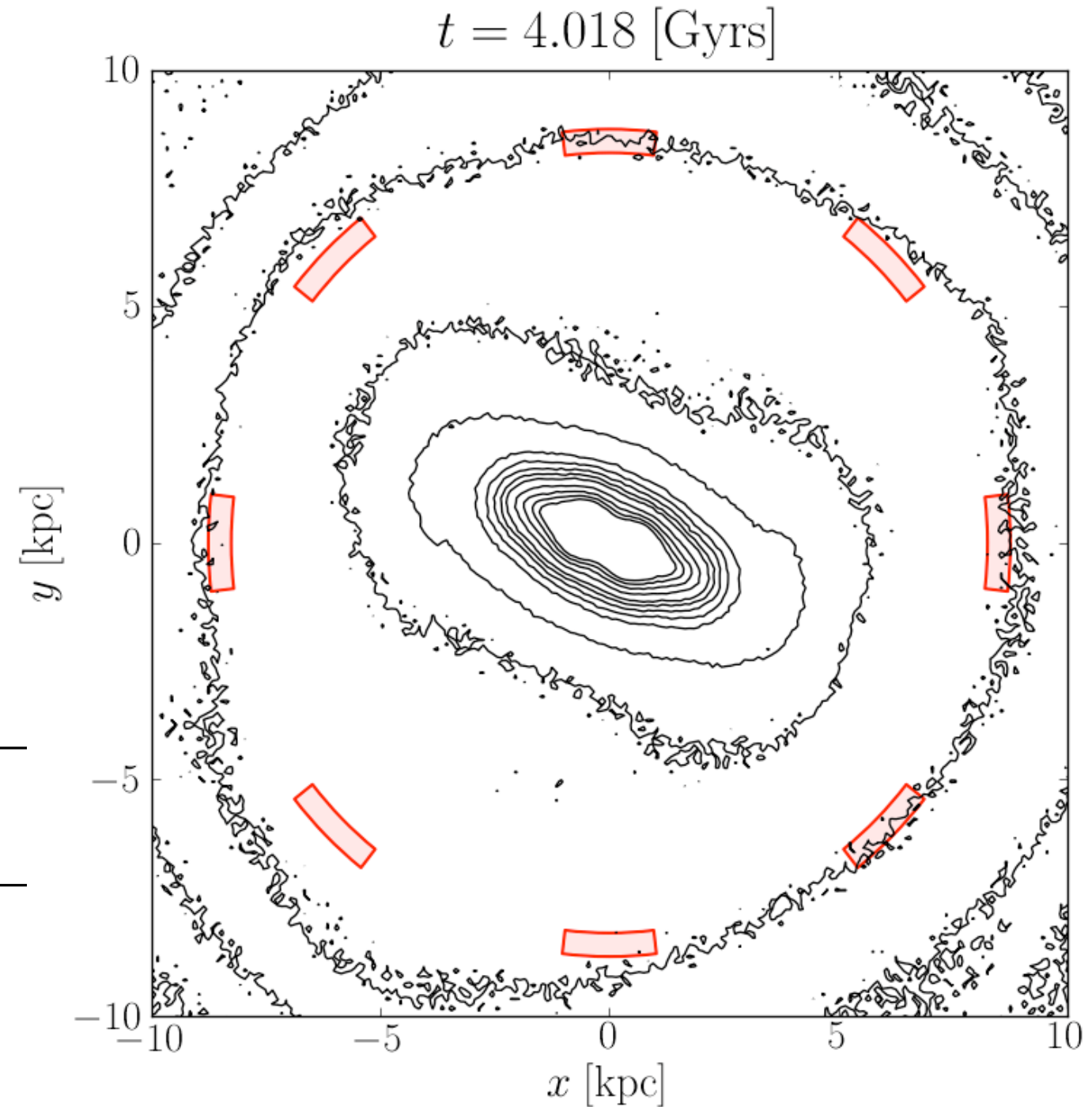
simulation

| | N (10^6) | M ($10^{10} M_{\odot}$) | ε (kpc) | $R_{1/2}$ (kpc) | $z_{1/2}$ (kpc) |
|-------|-----------------|--------------------------------|------------------------|--------------------|--------------------|
| Disc | 30 | 5.30 | 0.015 | 4.99 | 0.17 |
| Bulge | 0.5 | 0.83 | 0.012 | — | — |
| Halo | 15 | 45.40 | 0.045 | — | — |

Milky Way

| | M ($10^{10} M_{\odot}$) | $R_{1/2}$ (kpc) | $z_{1/2}$ (kpc) |
|------------|--------------------------------|--------------------|--------------------|
| Thin disc | 3.5–5.5 ^a | 3.35–9.24 | ~0.14–0.18 |
| Thick disc | — | 5.04–7.56 | 0.49–0.84 |
| Bulge | ~1 | — | — |
| Halo | ~40–200 | — | — |

^aTotal disc mass.



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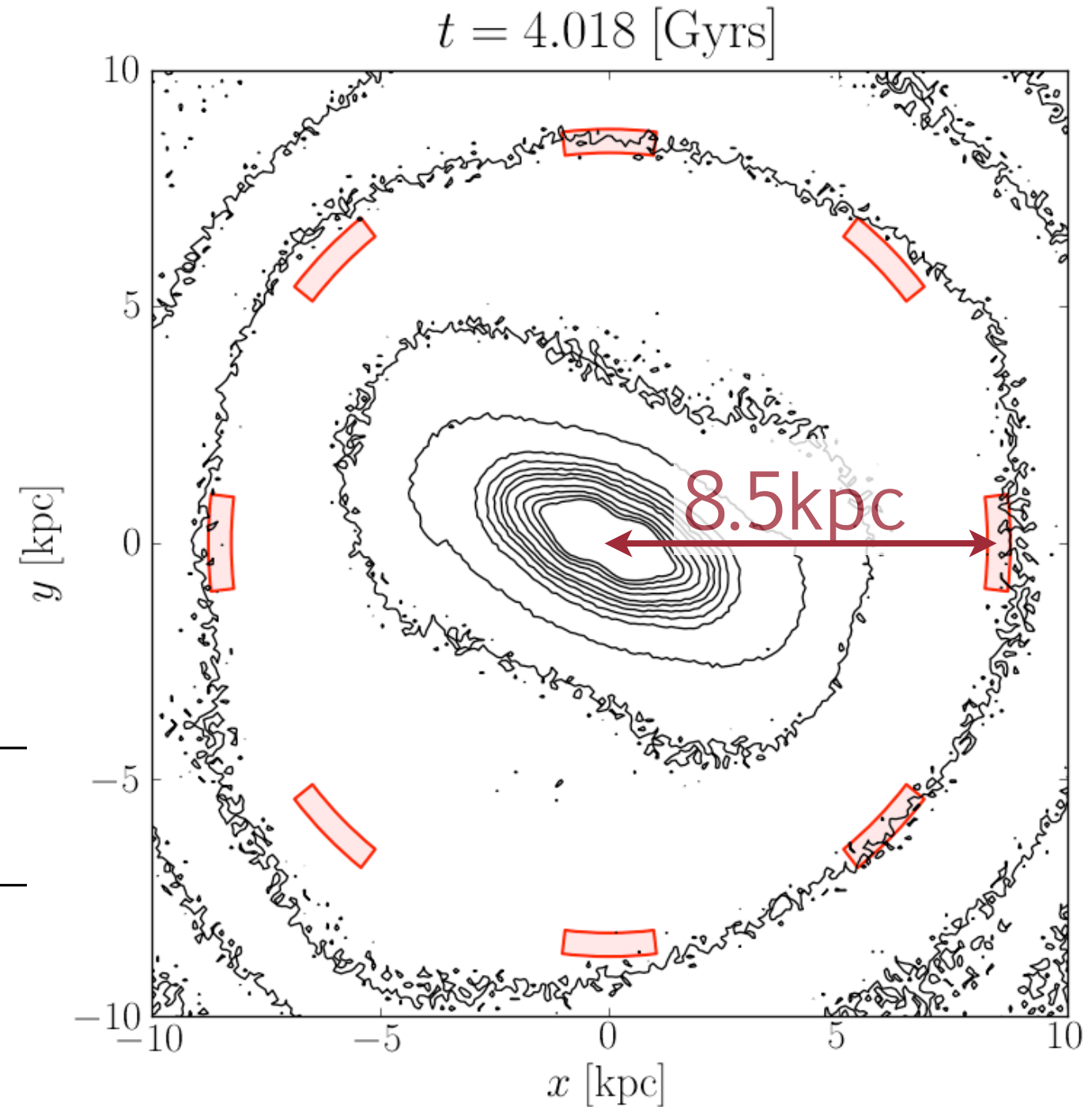
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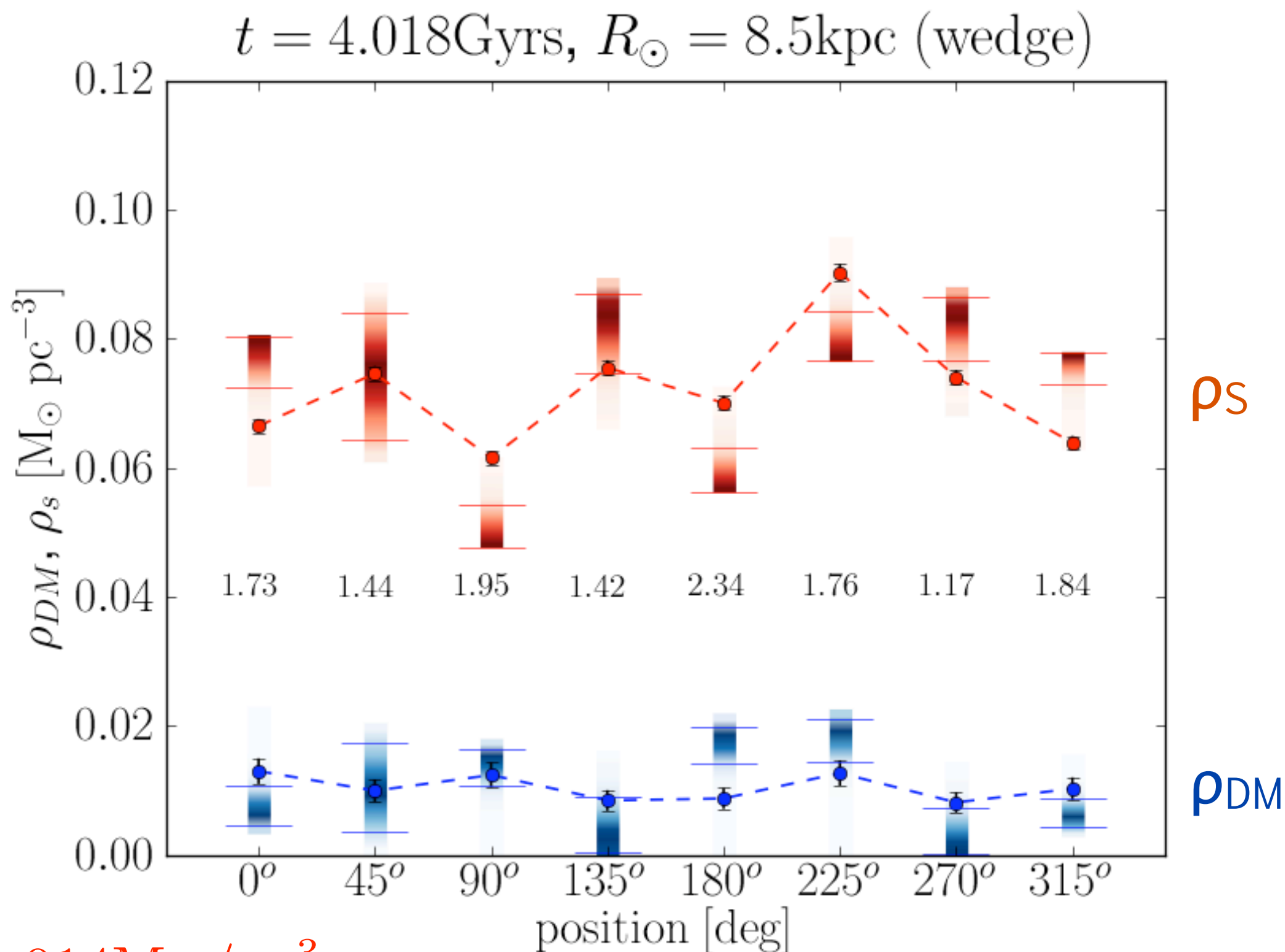
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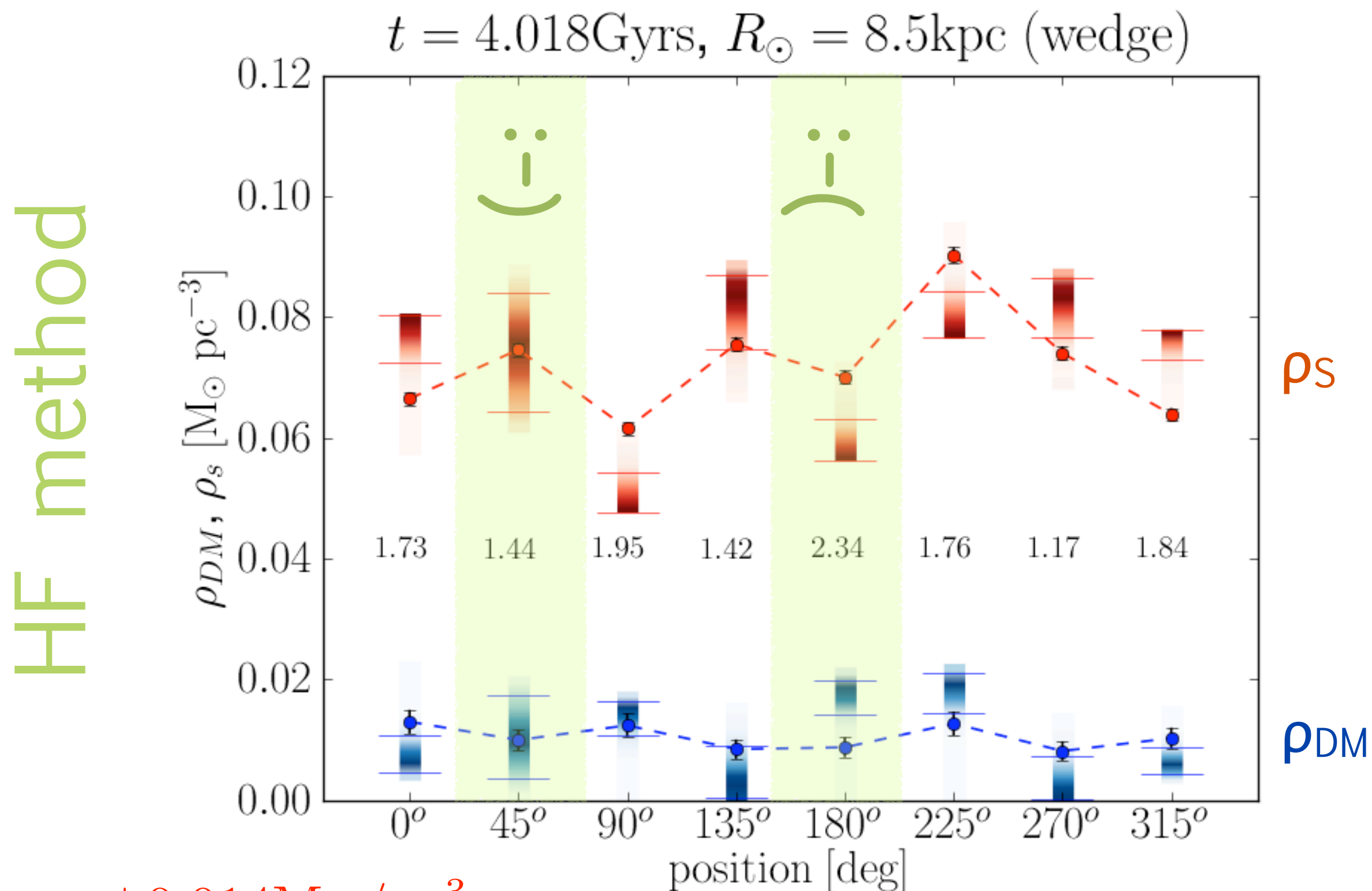
Evolved sim: HF vs MA method

HF method



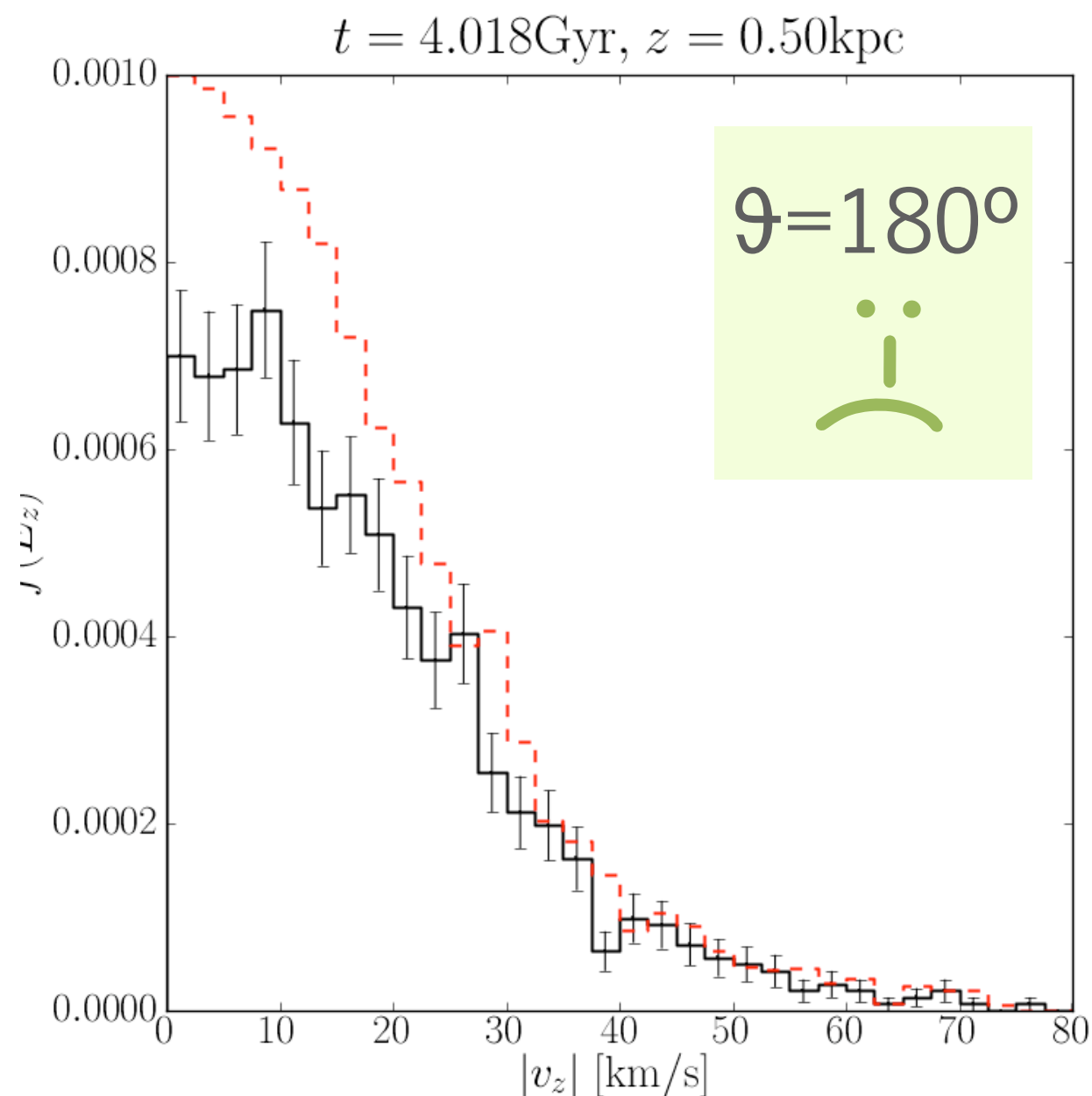
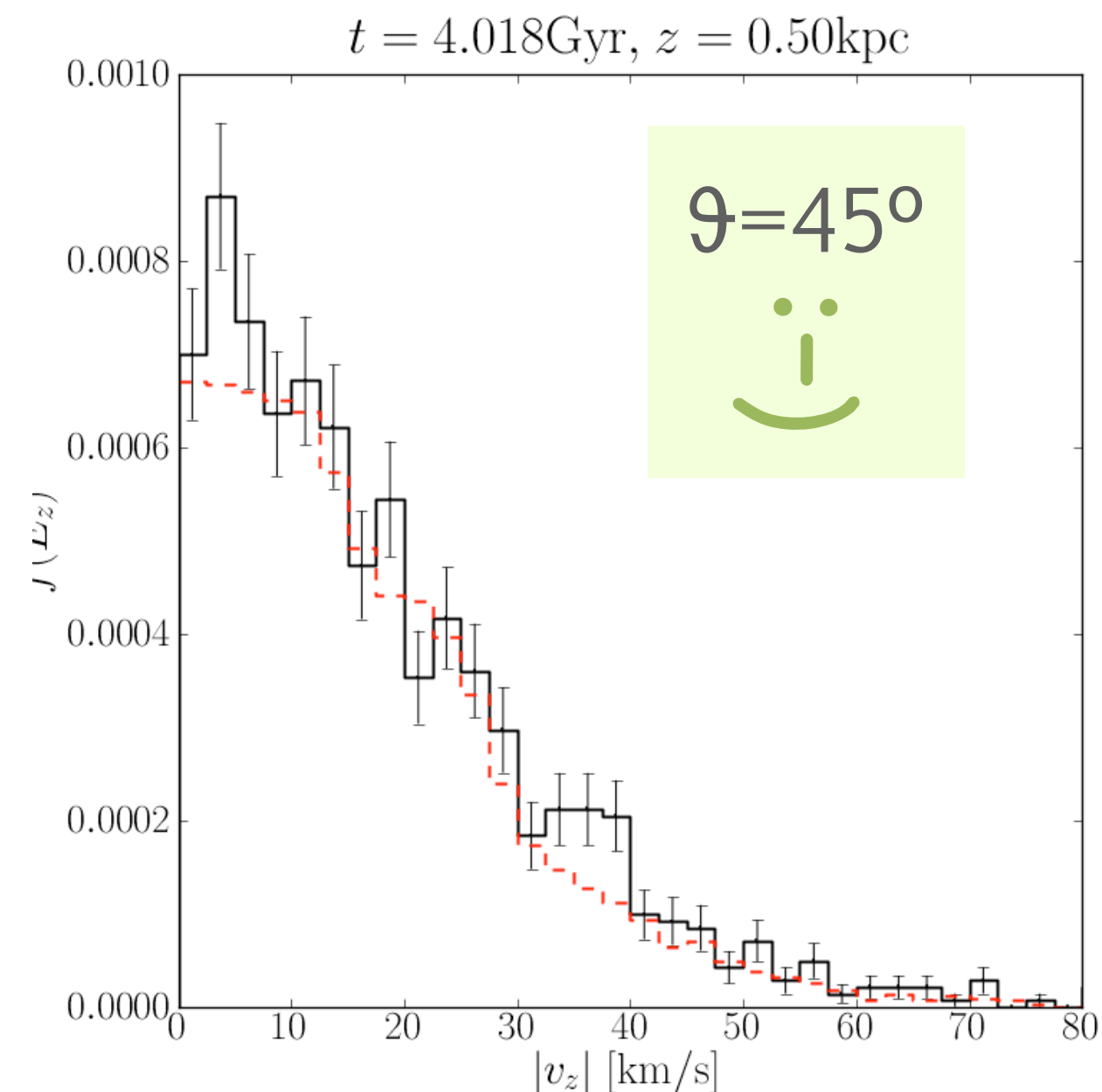
$$\Delta\rho_s = \pm 0.014 M_{\odot}/\text{pc}^3$$

Evolved sim: HF vs MA method



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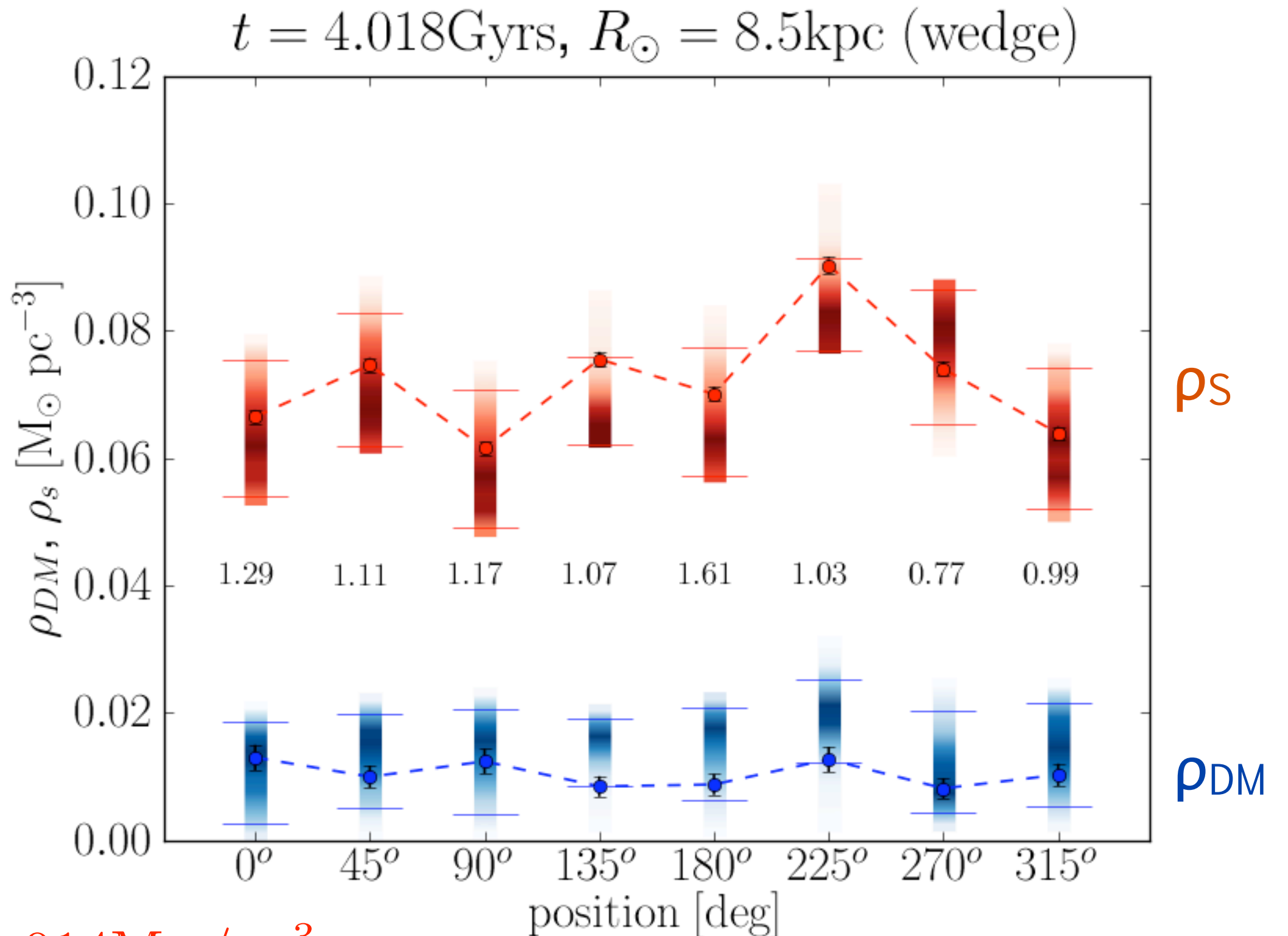
Evolved sim: distribution functions



Evolved simulation

Evolved sim: HF vs MA method

MA method



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Tracer population

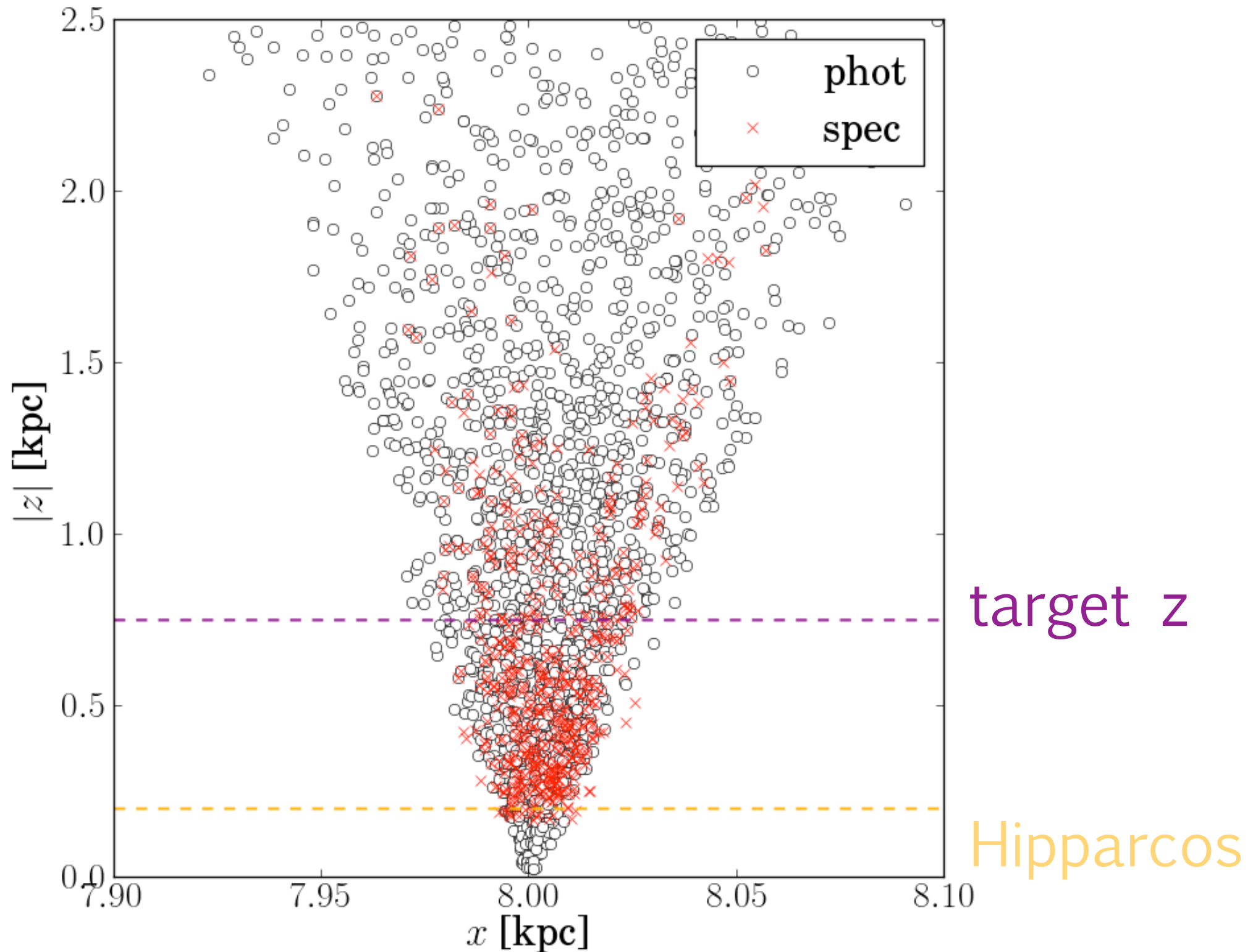
MUST BE:

- * in dynamical equilibrium with the Galactic potential.
- * common stars (to allow useful statistical precision in the result).
- * in a volume complete sample.
- * with reliable distances and vertical velocity available.

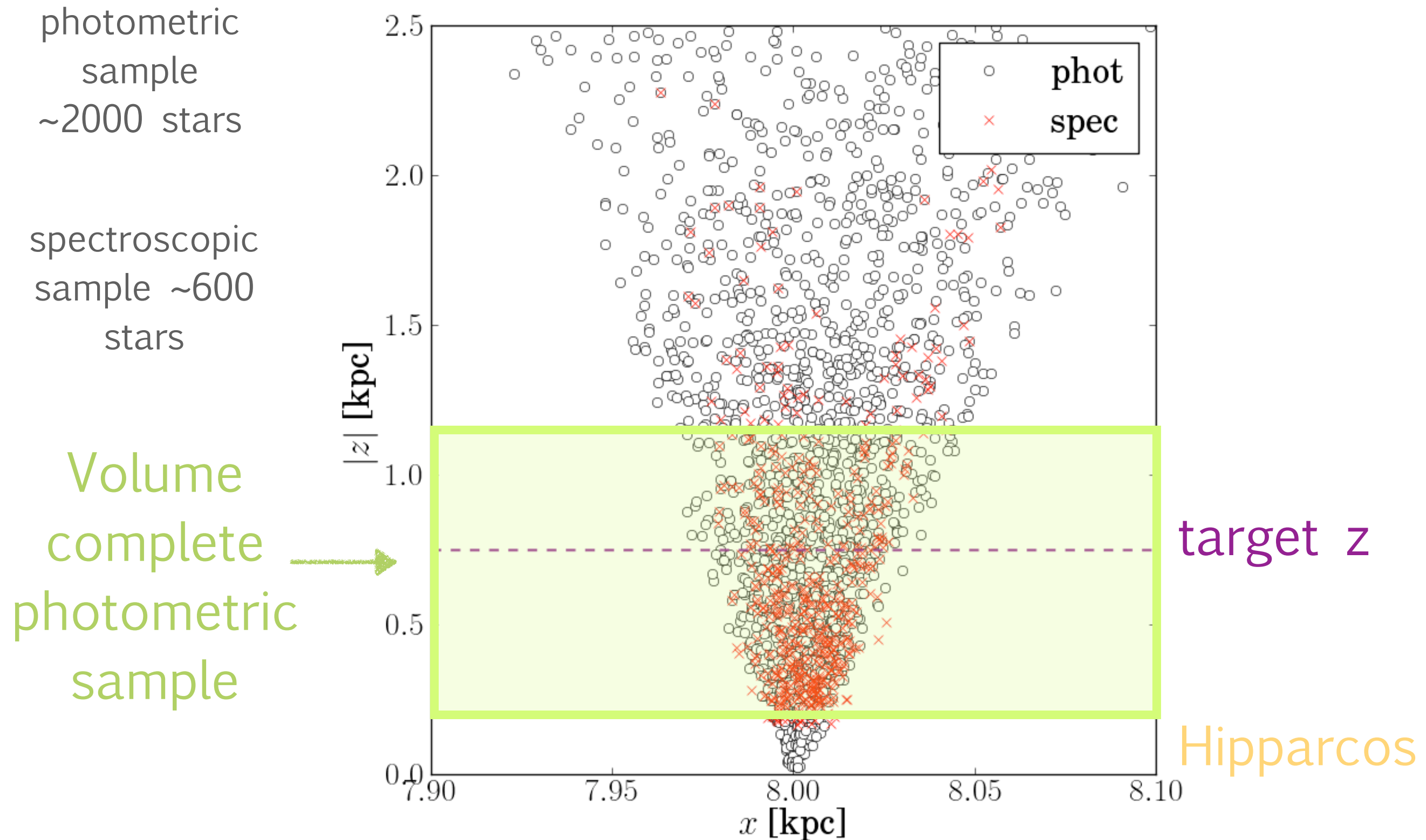
DATA: from Kuijken & Gilmore 1989 - K dwarfs

photometric
sample
~2000 stars

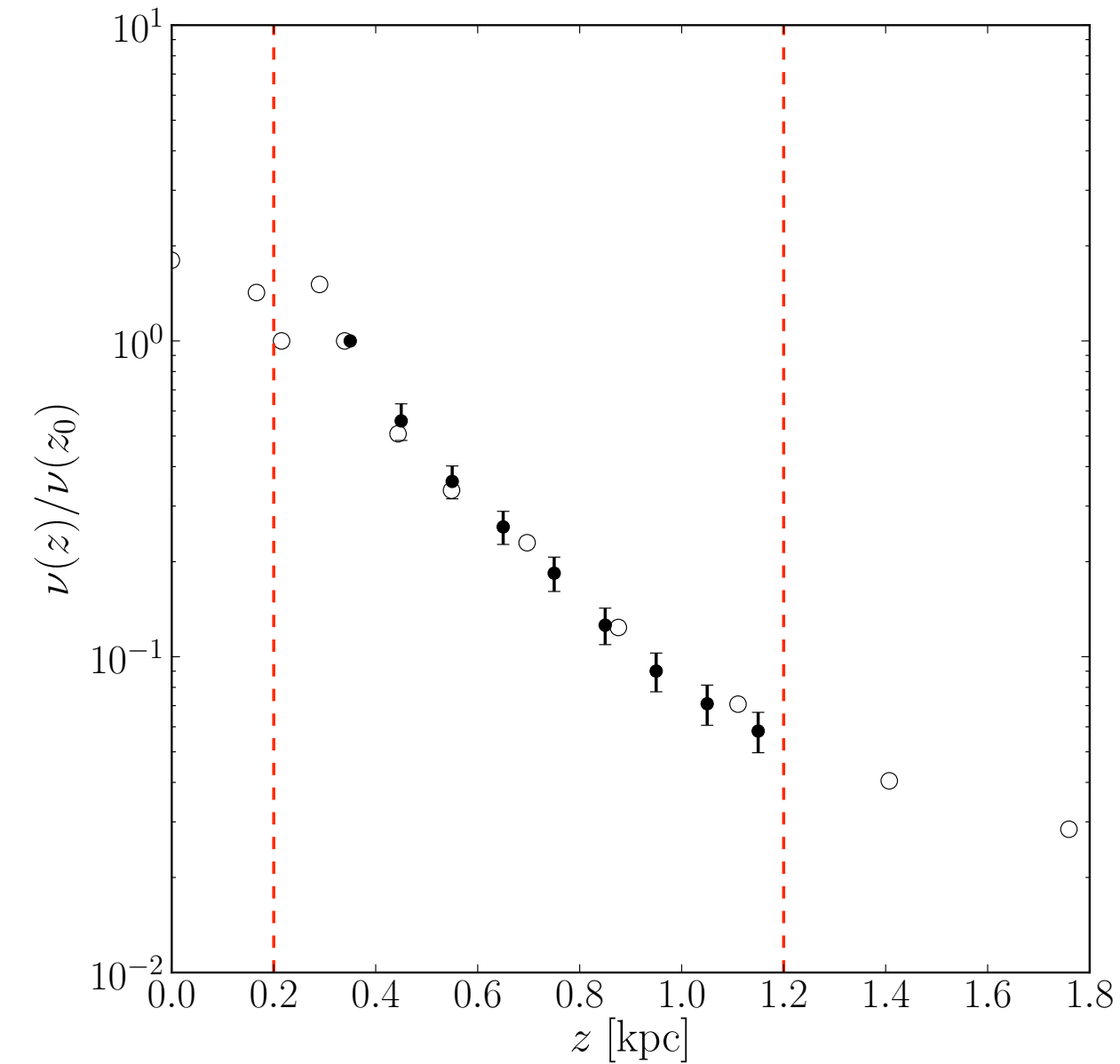
spectroscopic
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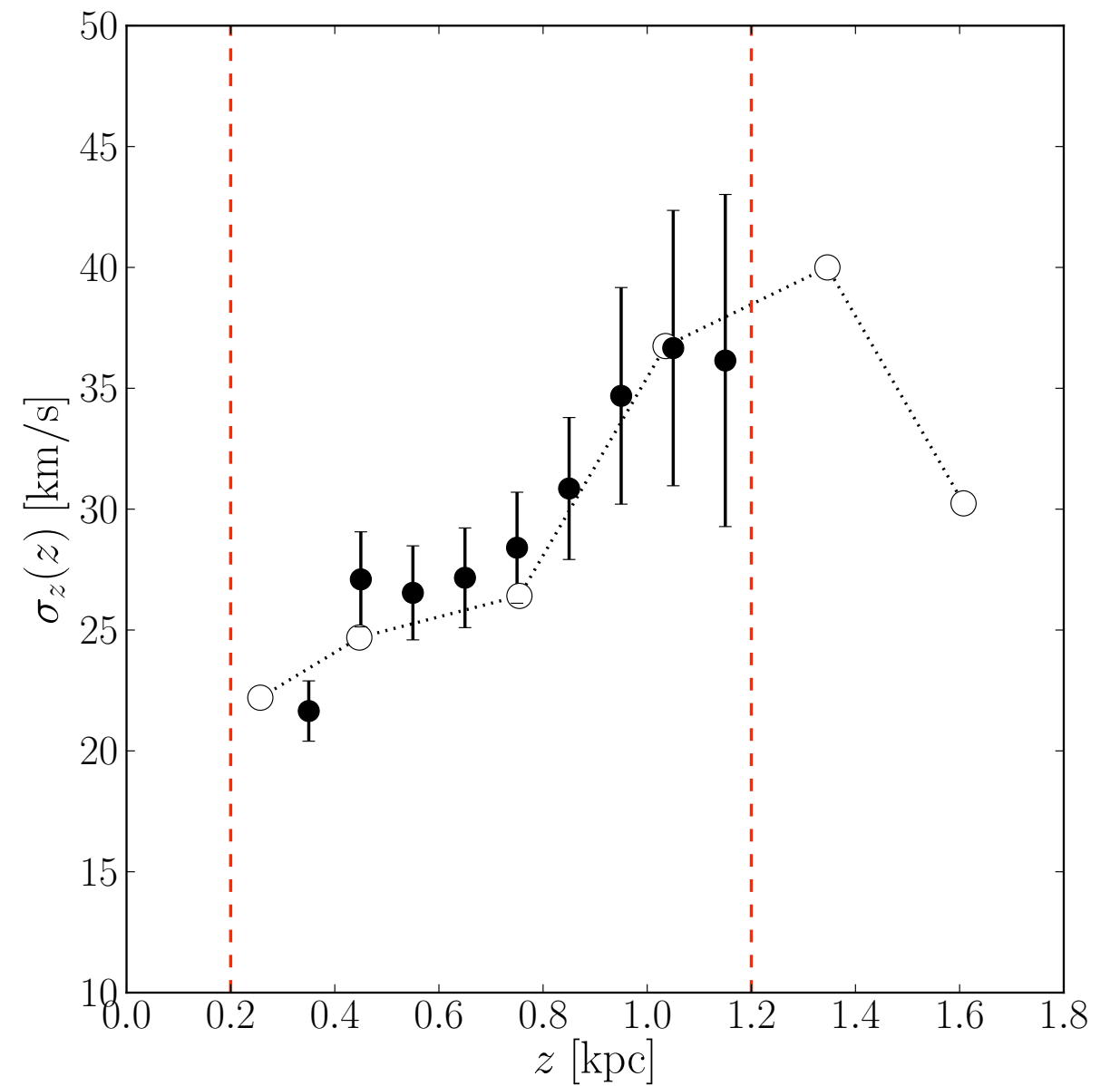
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Visible mass model by Flynn et al 2006

$$\rho_s(z) = \sum_i \nu_{i,0} \exp \left(-\frac{\Phi(z)}{v_{z,i}^2} \right)$$

density errors:
Stars: 10-20%;
Gas*: 50%

| Component | $\nu_{i,0}(0)$ [M \odot /pc 3] | $\overline{v_{z,i}^2}(0)$ [km/s] |
|-------------------|---|-------------------------------------|
| H $_2^*$ | 0.021 | 4.0 \pm 1.0 |
| HI(1)* | 0.016 | 7.0 \pm 1.0 |
| HI(2)* | 0.012 | 9.0 \pm 1.0 |
| Warm gas* | 0.0009 | 40.0 \pm 1.0 |
| Giants | 0.0006 | 20.0 \pm 2.0 |
| $M_V < 2.5$ | 0.0031 | 7.5 \pm 2.0 |
| $2.5 < M_V < 3.0$ | 0.0015 | 10.5 \pm 2.0 |
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| $4.0 < M_V < 5.0$ | 0.0022 | 18.0 \pm 2.0 |
| $5.0 < M_V < 8.0$ | 0.007 | 18.5 \pm 2.0 |
| $M_V > 8.0$ | 0.0135 | 18.5 \pm 2.0 |
| White dwarfs | 0.006 | 20.0 \pm 5.0 |
| Brown dwarfs | 0.002 | 20.0 \pm 5.0 |
| Thick disc | 0.0035 | 37.0 \pm 5.0 |
| Stellar halo | 0.0001 | 100.0 \pm 10.0 |

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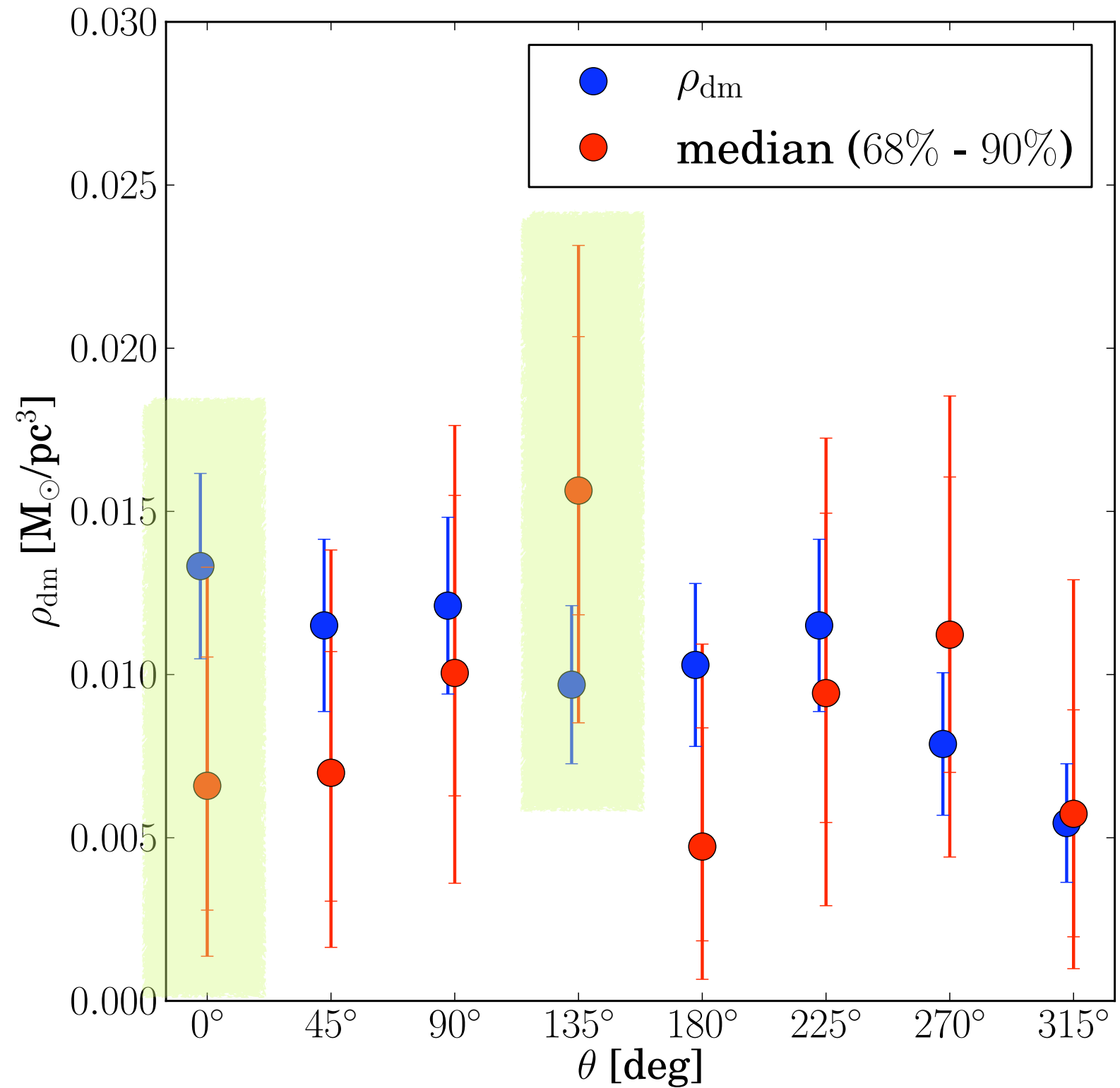
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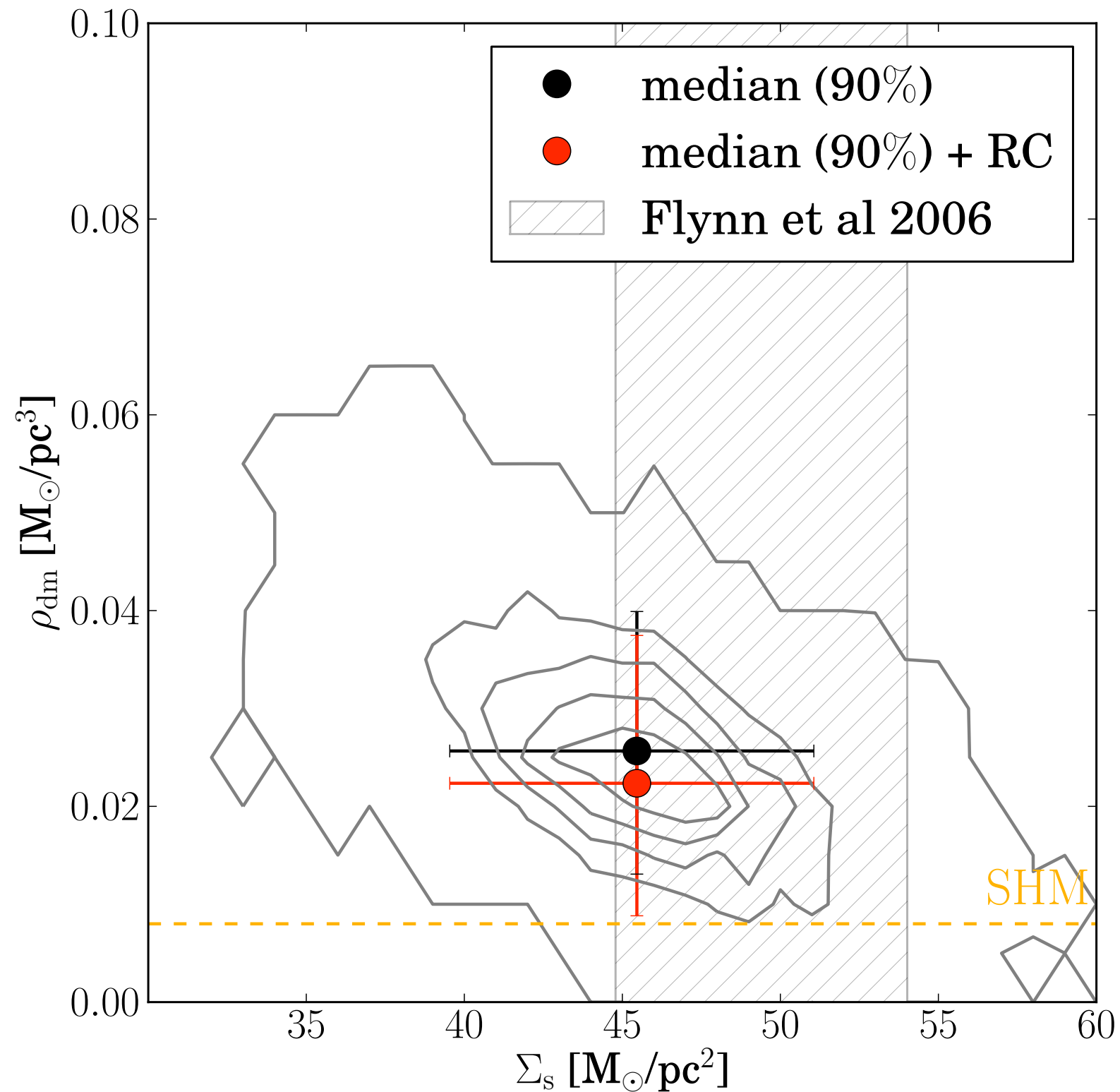
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| Component | $\nu_{i,0}(0)$ [M \odot /pc 3] | $\overline{v_{z,i}^2}(0)$ [km/s] |
|-------------------|---|-------------------------------------|
| H $_2^*$ | 0.021 | 4.0 \pm 1.0 |
| HI(1)* | 0.016 | 7.0 \pm 1.0 |
| HI(2)* | 0.012 | 9.0 \pm 1.0 |
| Warm gas* | 0.0009 | 40.0 \pm 1.0 |
| Giants | 0.0006 | 20.0 \pm 2.0 |
| $M_V < 2.5$ | 0.0031 | 7.5 \pm 2.0 |
| $2.5 < M_V < 3.0$ | 0.0015 | 10.5 \pm 2.0 |
| $3.0 < M_V < 4.0$ | 0.0020 | 14.0 \pm 2.0 |
| $4.0 < M_V < 5.0$ | 0.0022 | 18.0 \pm 2.0 |
| $5.0 < M_V < 8.0$ | 0.007 | 18.5 \pm 2.0 |
| $M_V > 8.0$ | 0.0135 | 18.5 \pm 2.0 |
| White dwarfs | 0.006 | 20.0 \pm 5.0 |
| Brown dwarfs | 0.002 | 20.0 \pm 5.0 |
| Thick disc | 0.0035 | 37.0 \pm 5.0 |
| Stellar halo | 0.0001 | 100.0 \pm 10.0 |

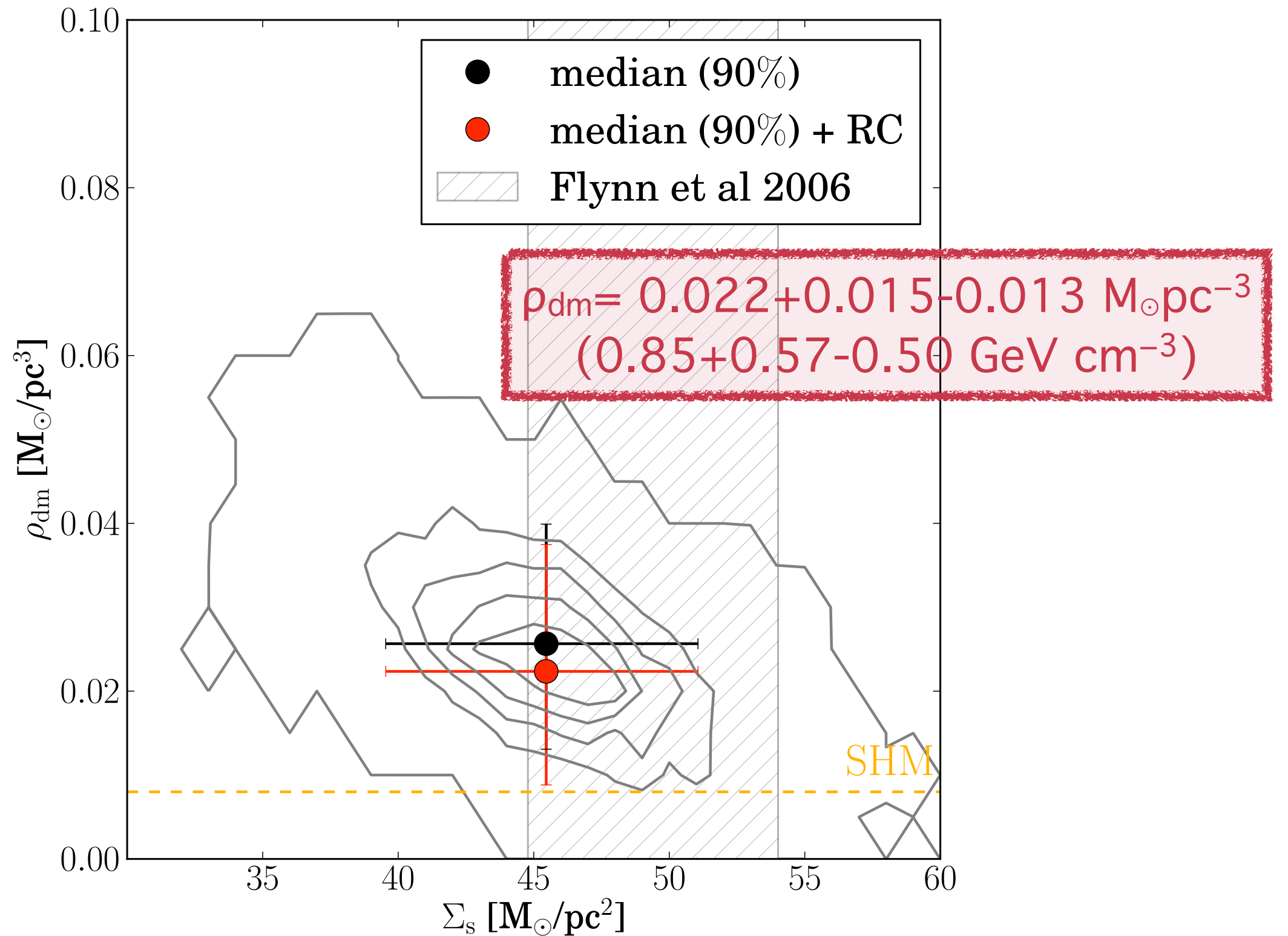
Test on the simulation: MA method



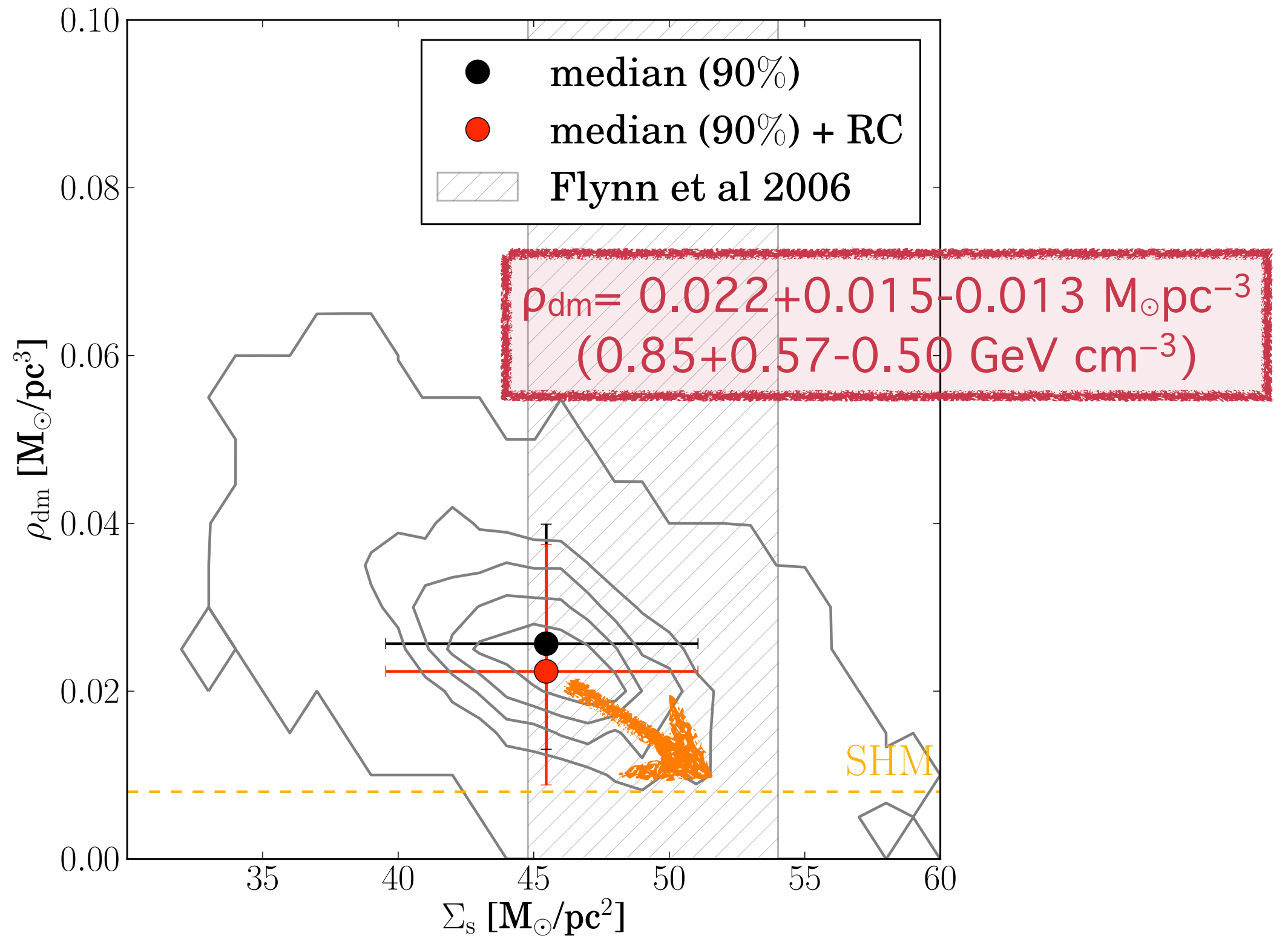
data from Kuijken & Gilmore 1989 - K dwarfs



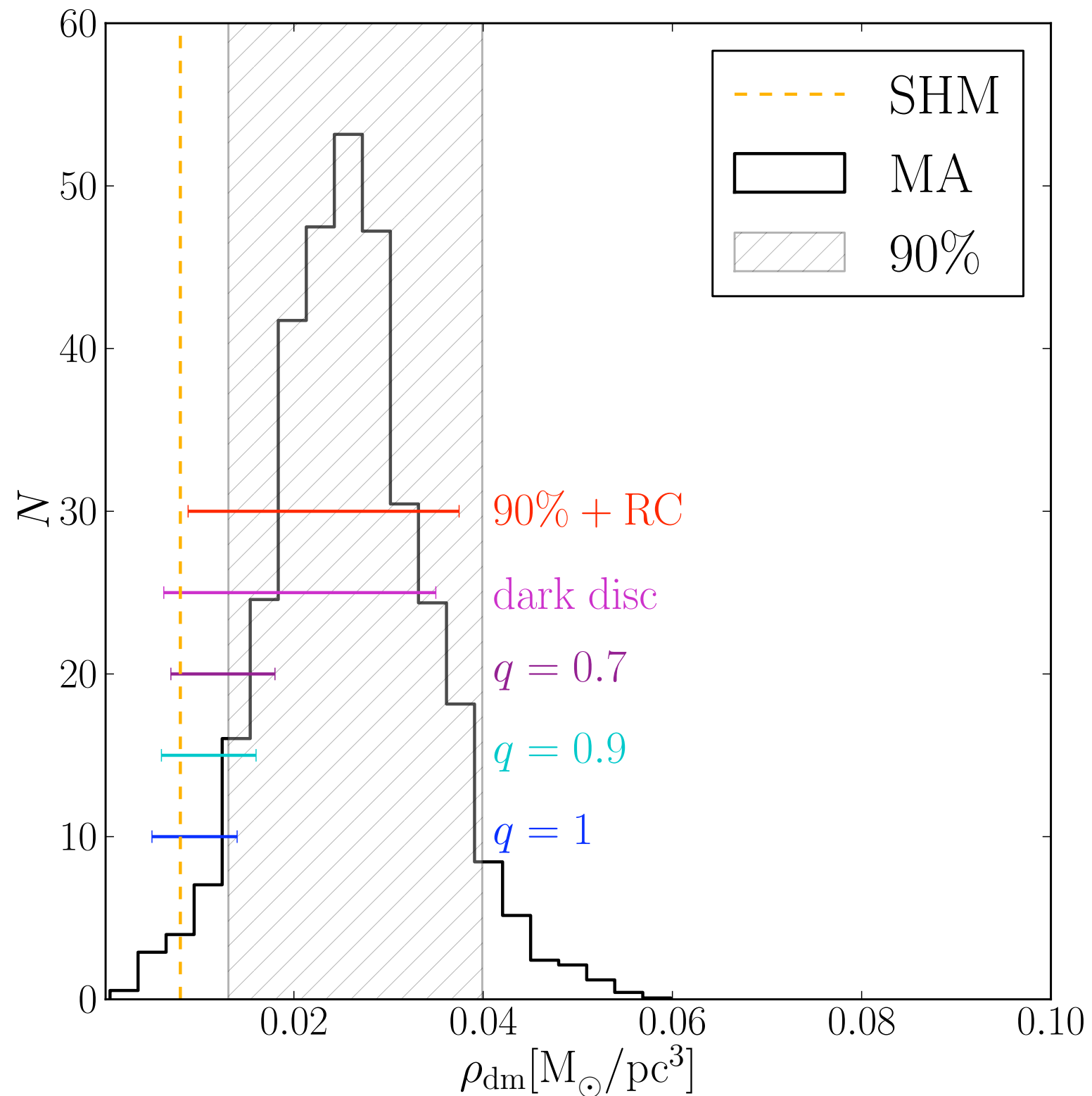
data from Kuijken & Gilmore 1989 - K dwarfs



data from Kuijken & Gilmore 1989 - K dwarfs



data from Kuijken & Gilmore 1989 - K dwarfs



Conclusion

* We present a new method to measure ρ_{dm} from the vertical kinematics local tracers. It relies on a minimal set of assumptions and

- uses a MCMC to marginalise over the uncertainties
- does not require any prior on the MW rotation curve
- does not require any assumption on the tracers' distribution function

* We use hi-res simulations as a mock data set to test our method.

* We obtain a new measurement of the local dark matter density: $\rho_{\text{dm}} = 0.022 \pm 0.015 \text{--} 0.013 \text{ M}_{\odot} \text{pc}^{-3}$ ($0.85 \pm 0.57 \text{--} 0.50 \text{ GeV cm}^{-3}$).

* Our median value of the local dark matter density is larger at 90% confidence than the Standard Halo Model value of $\rho^{\text{SHM}}_{\text{dm}} = 0.008 \text{ M}_{\odot} \text{pc}^{-3}$ (0.30 GeV cm^{-3}). If confirmed by future data (GAIA), it has interesting implications:

- for direct detection experiments: it implies a larger flux of dark matter particles and therefore a greater chance of detection.
- it suggests that the halo of our Galaxy is oblate and/or that we have a disc of dark matter.